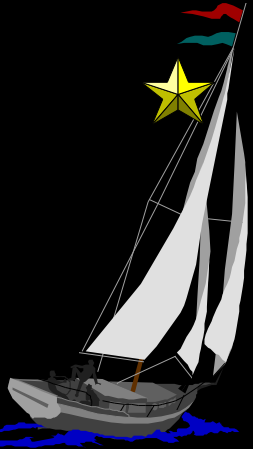


# The Properties of Winds from Massive Stars



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**Asif ud-Doula**

**Rich Townsend**

**Sapporo, Japan**

**IAR, Argentina**

**Swarthmore Coll.**

**Penn State**

**U. Wisconsin**

**& many more...**

# Key properties of massive-star winds

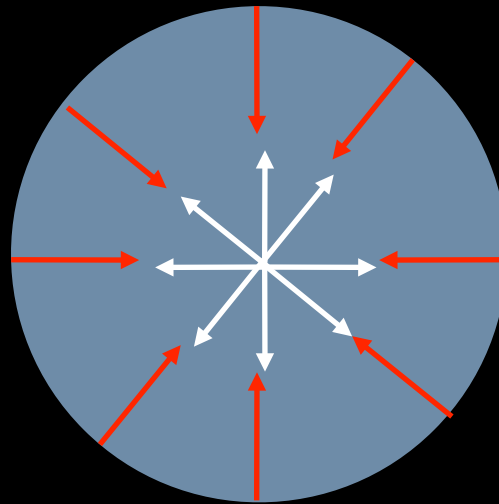
- $\dot{M}_{\text{dot}} \sim 10^{-10} - 10^{-4} M_{\text{sun}}/\text{yr}$ ;
  - $V_{\text{inf}} \sim 3 V_{\text{esc}} \sim 1000 - 3000 \text{ km/s}$
- sometimes (e.g. Be stars) rotationally distorted
  - **polar** wind & equatorial “**decretion disk**”
- highly **clumped** & moderately “**porous**”
  - filling factor  $f_{\text{vol}} \sim 0.1$
  - mfp  $\sim$  “porosity length”  $h \sim \ell / f_{\text{vol}} \sim 1 - 10\% R_*$

# Radiative force vs. gravity

Radiative  
Force

Gravitational  
Force

$$g_{rad} = \int_0^{\infty} d\nu \frac{\kappa_{\nu} F_{\nu}}{c}$$



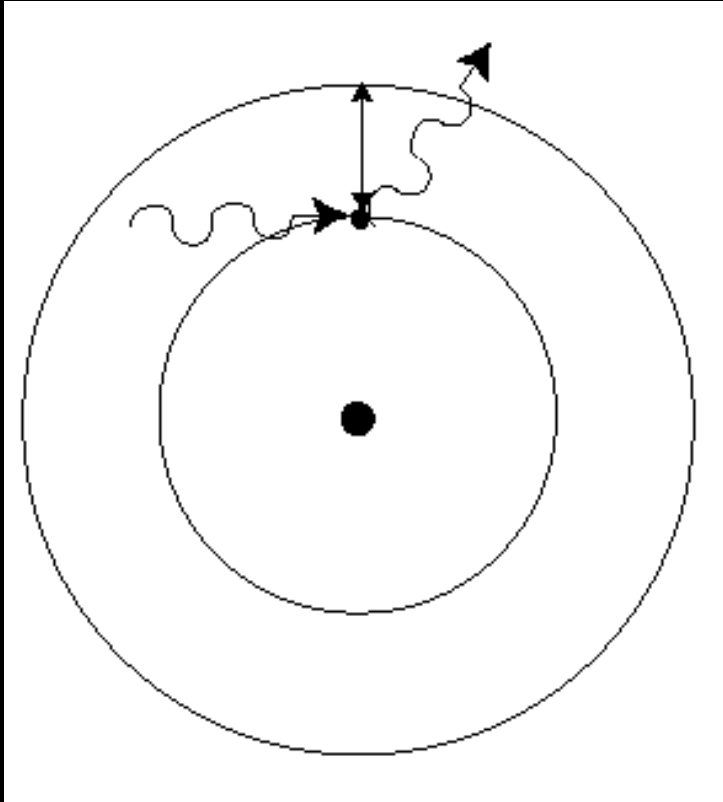
$$\frac{GM}{r^2}$$

$$\Gamma_e \equiv \frac{g_e}{g} = \frac{\kappa_e L / 4\pi r^2 c}{GM / r^2} = \frac{\kappa_e L}{4\pi GM c}$$

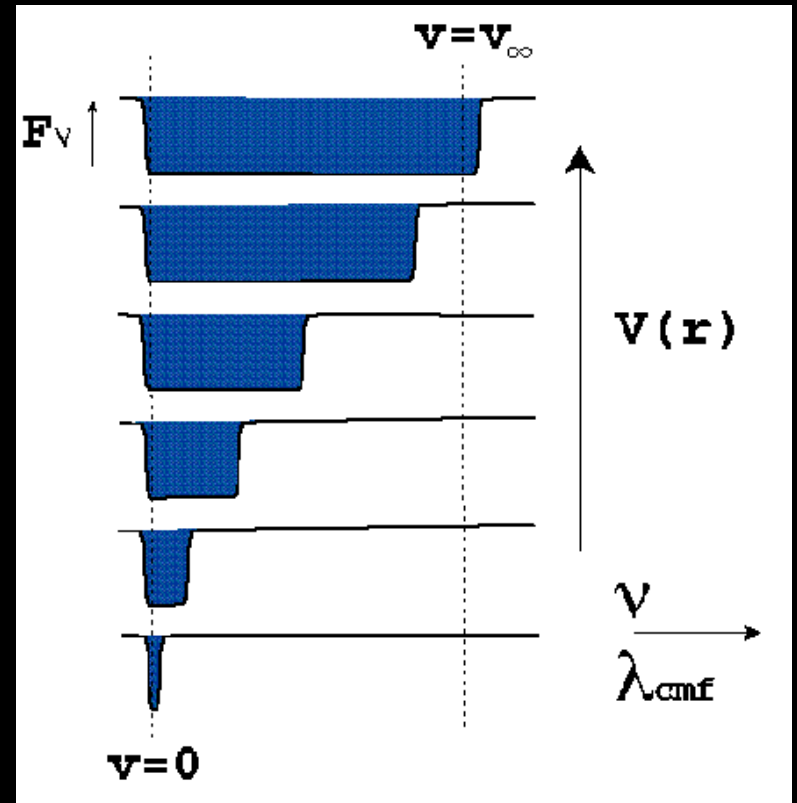
$$\Gamma \approx 2 \times 10^{-5} \frac{L / L_{\odot}}{M / M_{\odot}} \frac{\overline{\kappa_F}}{\kappa_e}$$

# Driving by **Line-Opacity**

Optically **thin**



Optically **thick**



$$\Gamma_{thin} \sim Q\Gamma_e \sim 1000\Gamma_e$$

$$\Gamma_{thick} \sim \frac{Q\Gamma_e}{\tau} \sim \frac{1}{\rho} \frac{dv}{dr}$$

# CAK model of steady-state wind

$0 < \alpha < 1$   
CAK ensemble of  
thick & thin lines

Equation of motion: 
$$v v' \approx -\frac{GM(1-\Gamma)}{r^2} + \frac{\bar{Q}L}{r^2} \left( \frac{r^2 v v'}{\dot{M} \bar{Q}} \right)^\alpha$$

inertia  $\approx$  gravity  $\approx$  CAK line-force

$g_{\text{CAK}} \approx$  gravity

Mass loss rate

$$\dot{M} \approx \frac{L}{c^2} \left( \frac{\bar{Q}\Gamma}{1-\Gamma} \right)^{\frac{1}{\alpha}-1}$$

inertia  $\approx$  gravity

Velocity law

$$v(r) \approx v_\infty (1 - R_* / r)^\beta \quad \beta \approx 0.8$$

$\sim v_{\text{esc}}$

**Wind-Momentum  
Luminosity law**

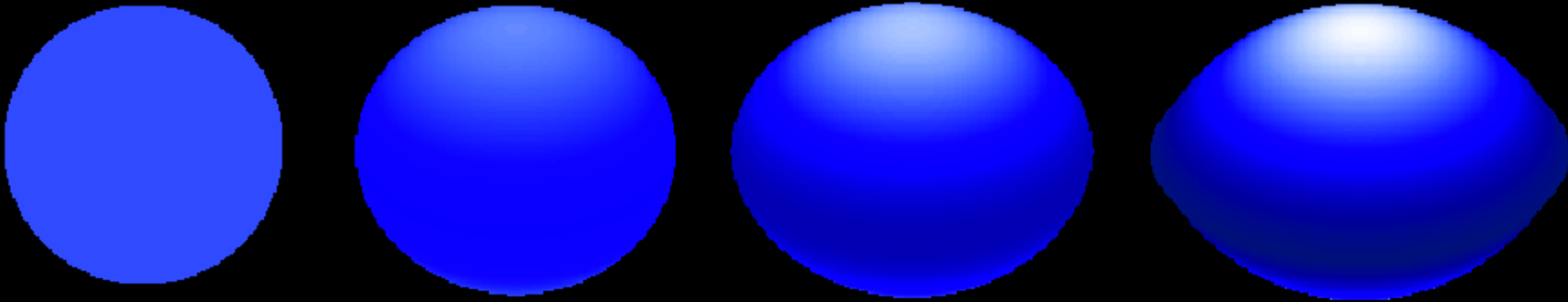
$$\dot{M} v_\infty \propto L^{\frac{1}{\alpha}}$$

$\alpha \approx 0.6$

**How is such a wind  
affected by (rapid)  
stellar rotation?**

# Gravity Darkening

**increasing stellar rotation** 



# Effect of gravity darkening on line-driven mass flux

Recall:

$$\dot{m}(\theta) \sim \frac{F(\theta)^{1/\alpha}}{g_{\text{eff}}(\theta)^{1/\alpha-1}} \sim \frac{F^2(\theta)}{g_{\text{eff}}(\theta)} \quad \text{e.g., for } \alpha = 1/2$$

w/o gravity darkening,  
if  $F(\theta) = \text{const.}$        $\dot{m}(\theta) \sim \frac{1}{g_{\text{eff}}(\theta)}$       highest at **equator**

w/ gravity darkening,  
if  $F(\theta) \sim g_{\text{eff}}(\theta)$        $\dot{m}(\theta) \sim F(\theta)$       highest at **pole**

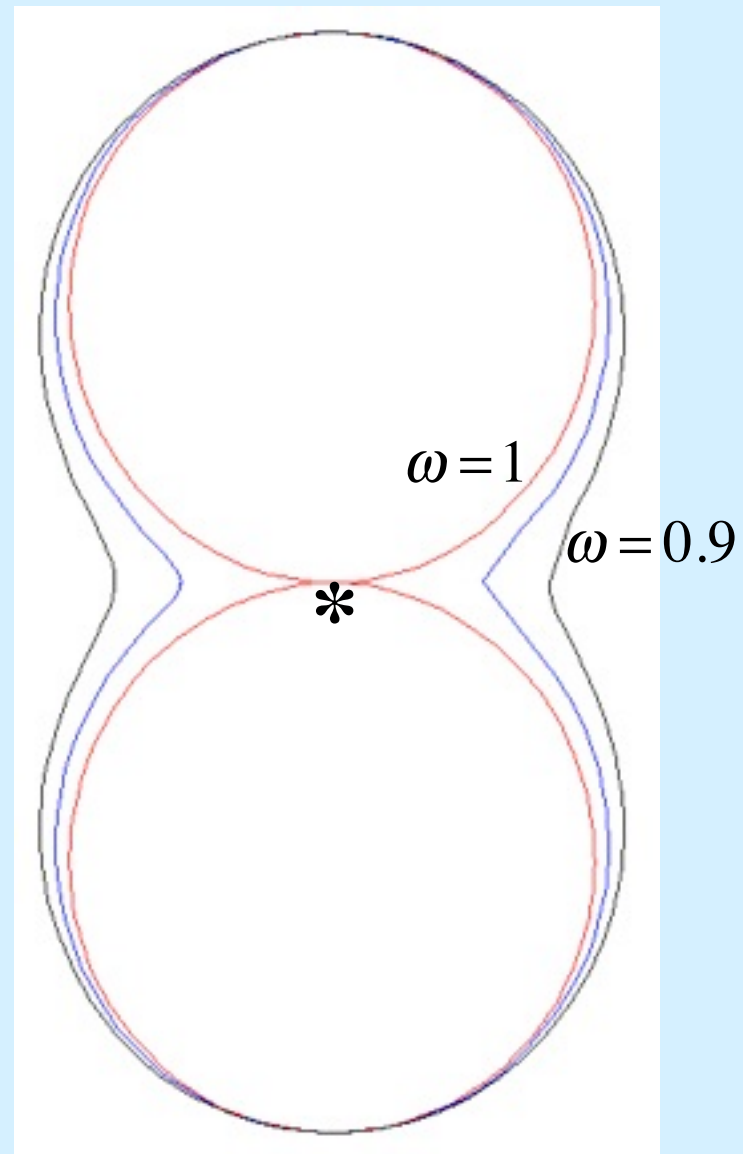


# Effect of rotation on flow speed

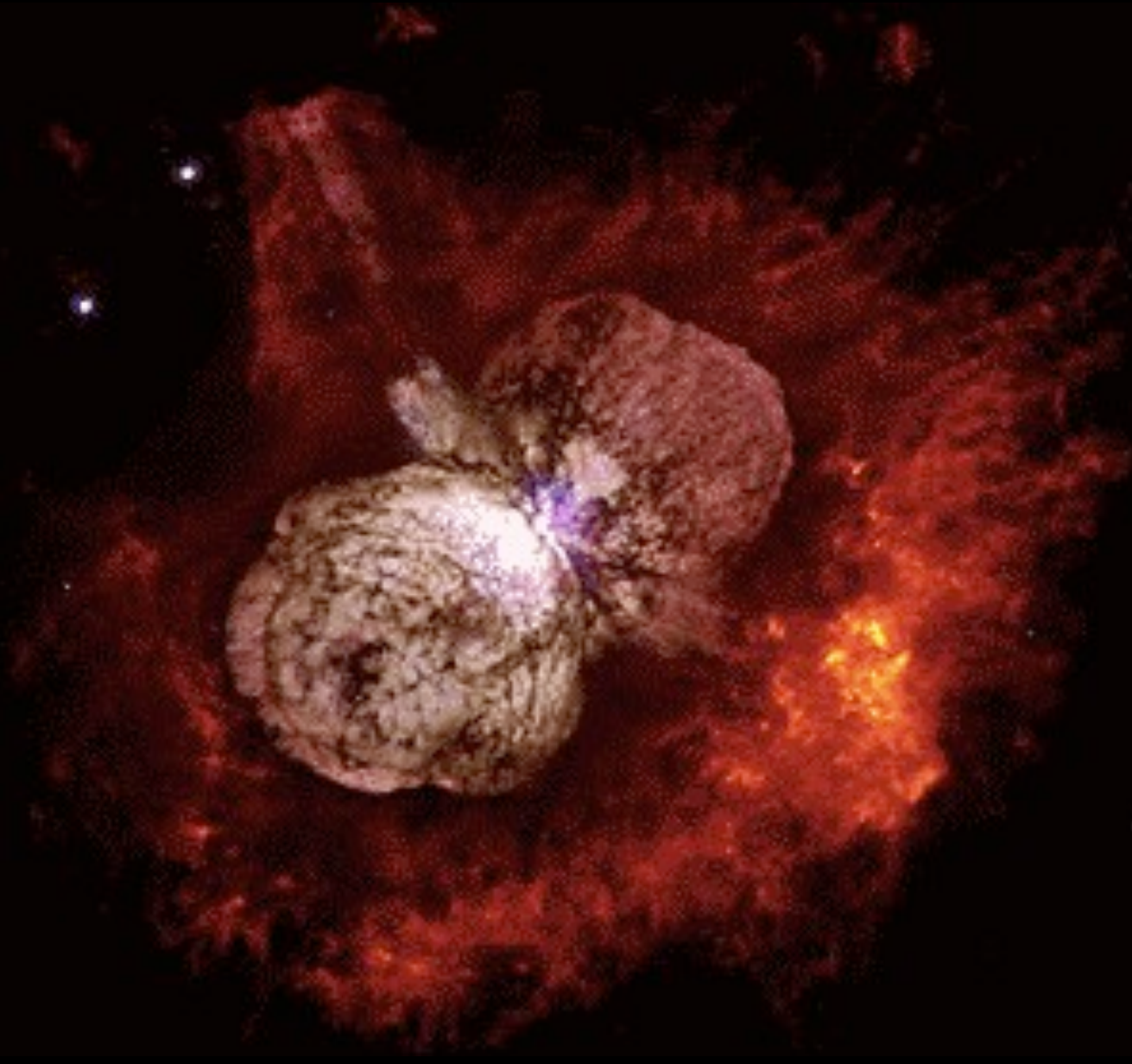
$$V_{\infty}(\theta) \sim V_{eff}(\theta) \sim \sqrt{g_{eff}(\theta)}$$

$$g_{eff}(\theta) \sim 1 - \omega^2 \sin^2 \theta$$

$$\omega \equiv \Omega / \Omega_{crit}$$



# Eta Carinae

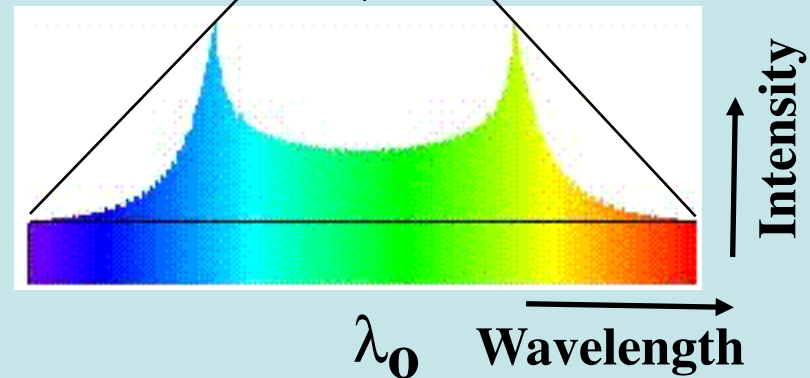


# Be stars

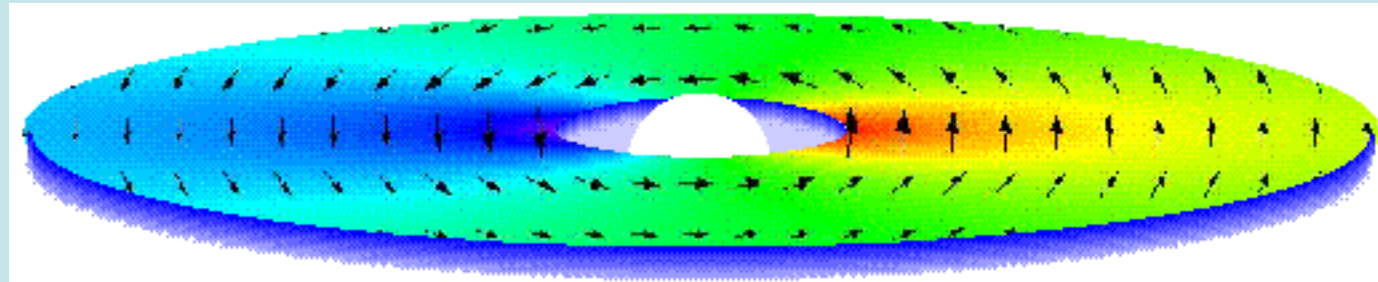
- Hot, bright, & **rapidly rotating** stars of mass  $\sim 3\text{-}10 M_{\text{sun}}$
- The “**e**” stands for **e**mission lines in the star’s spectrum



- Emission intensity split into **blue** and **red** peaks



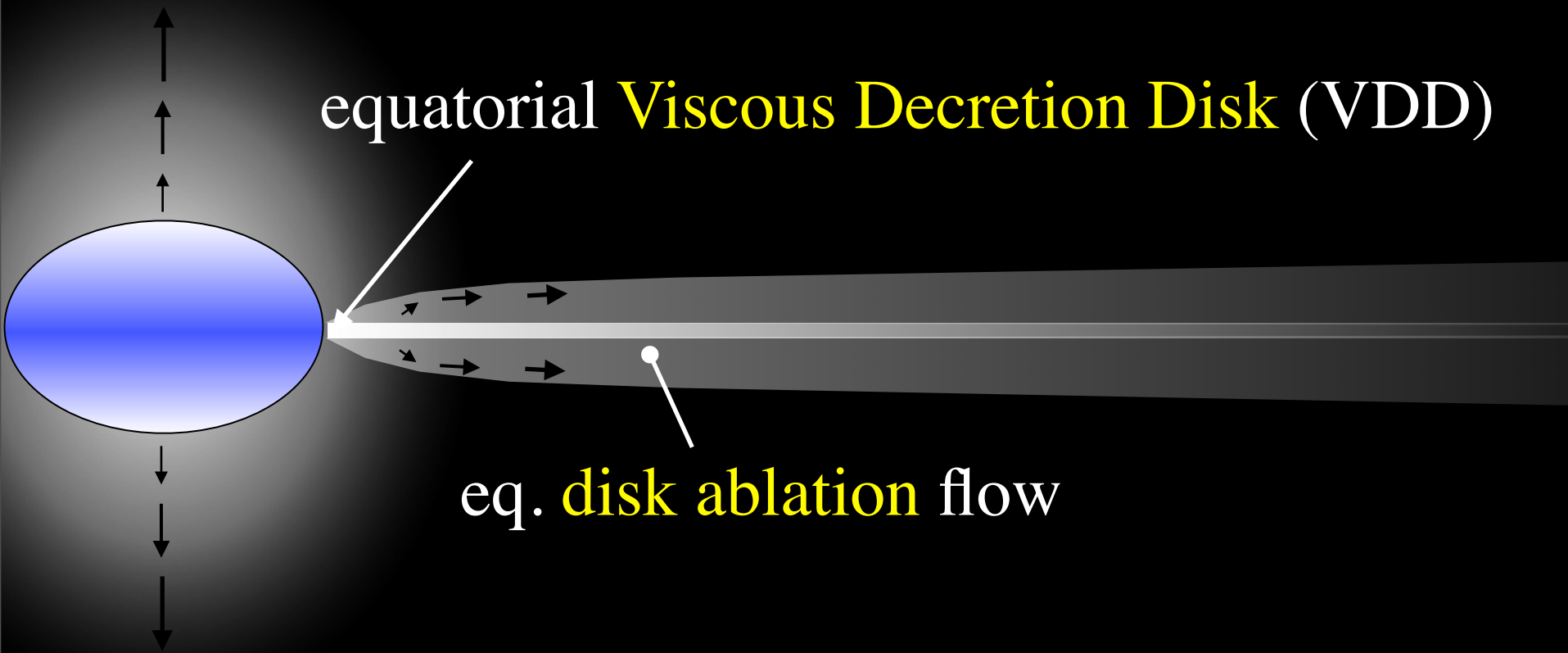
- From Doppler shift of gas moving **toward** and **away** from the observer .



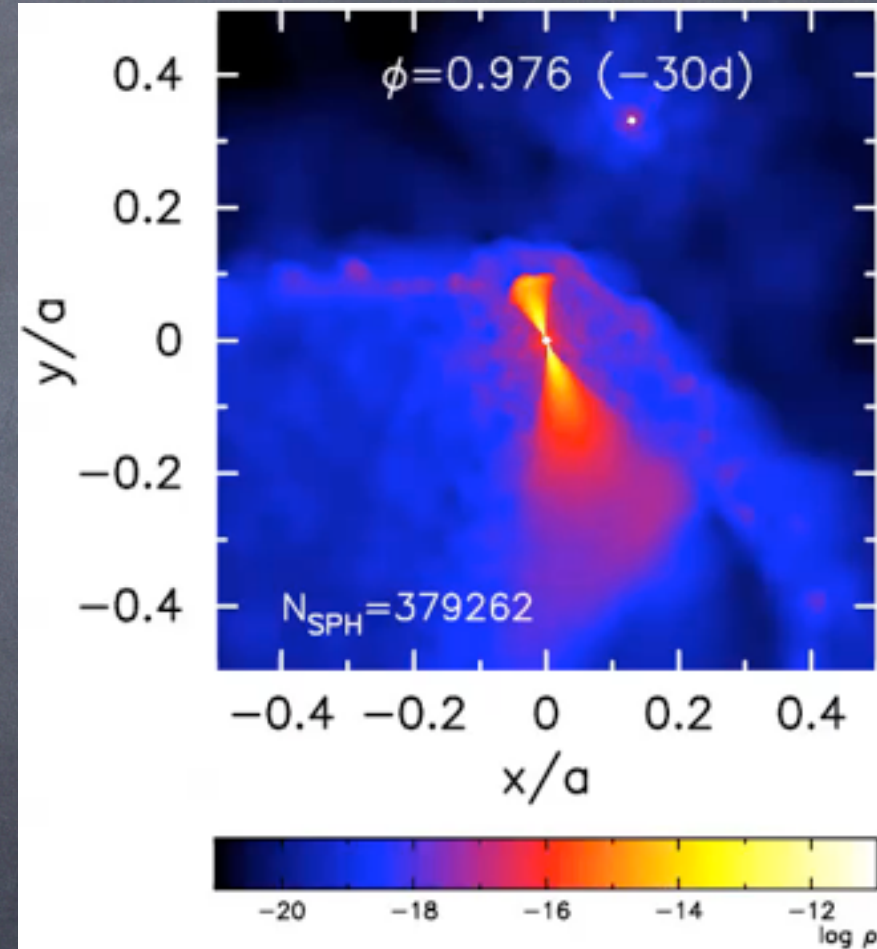
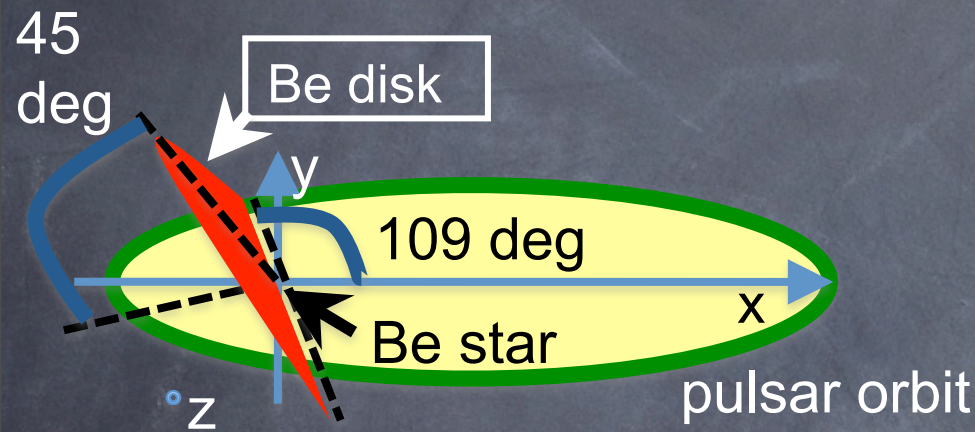
- Indicates a **disk of gas** orbits the star.

# 3 components of Be star circumstellar gas

gravity brightened poles  
drive denser **polar wind**



B1259-63 = misaligned Be-pulsar system.



see talk by A. Okazaki

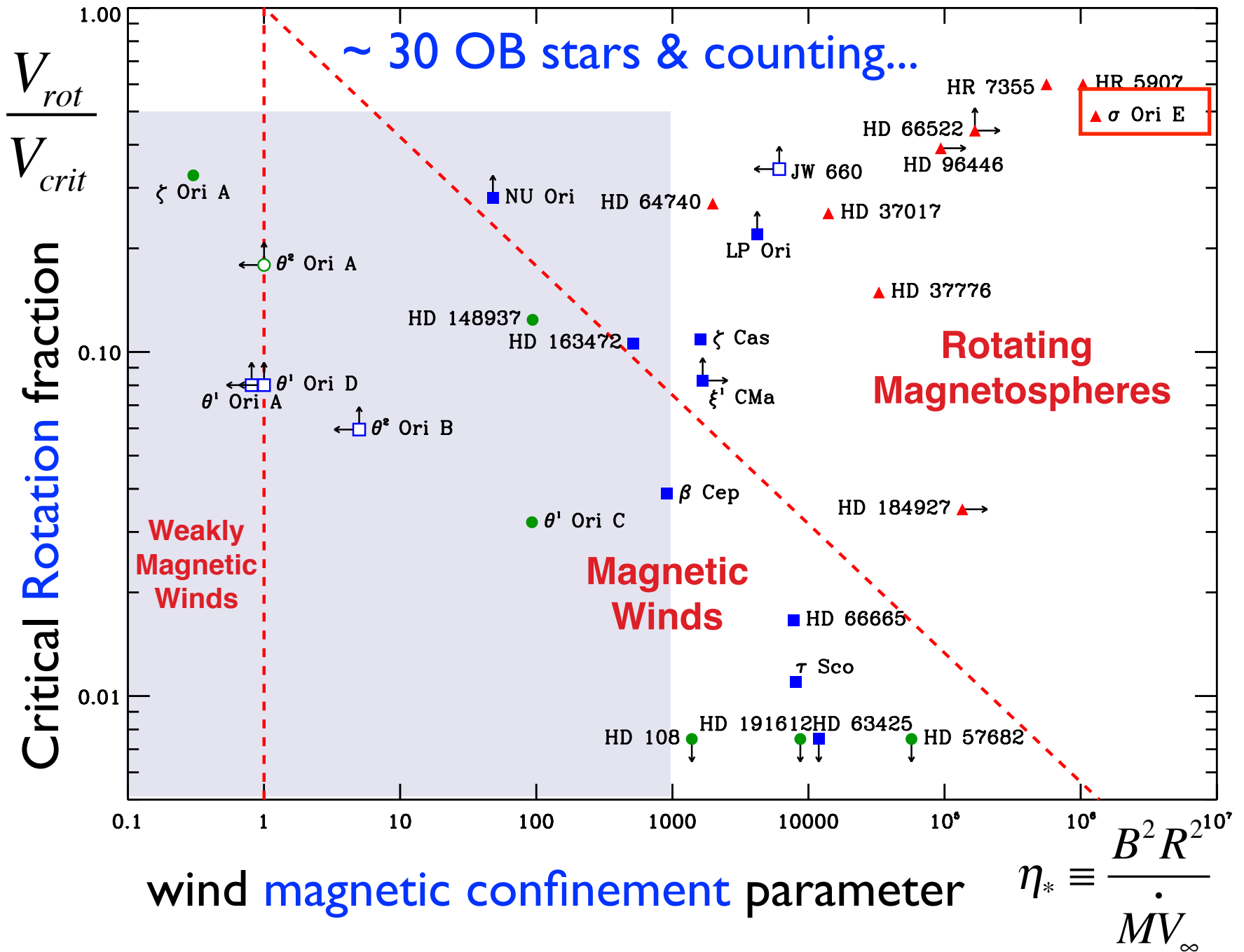
# MiMeS

## Magnetism in Massive Stars

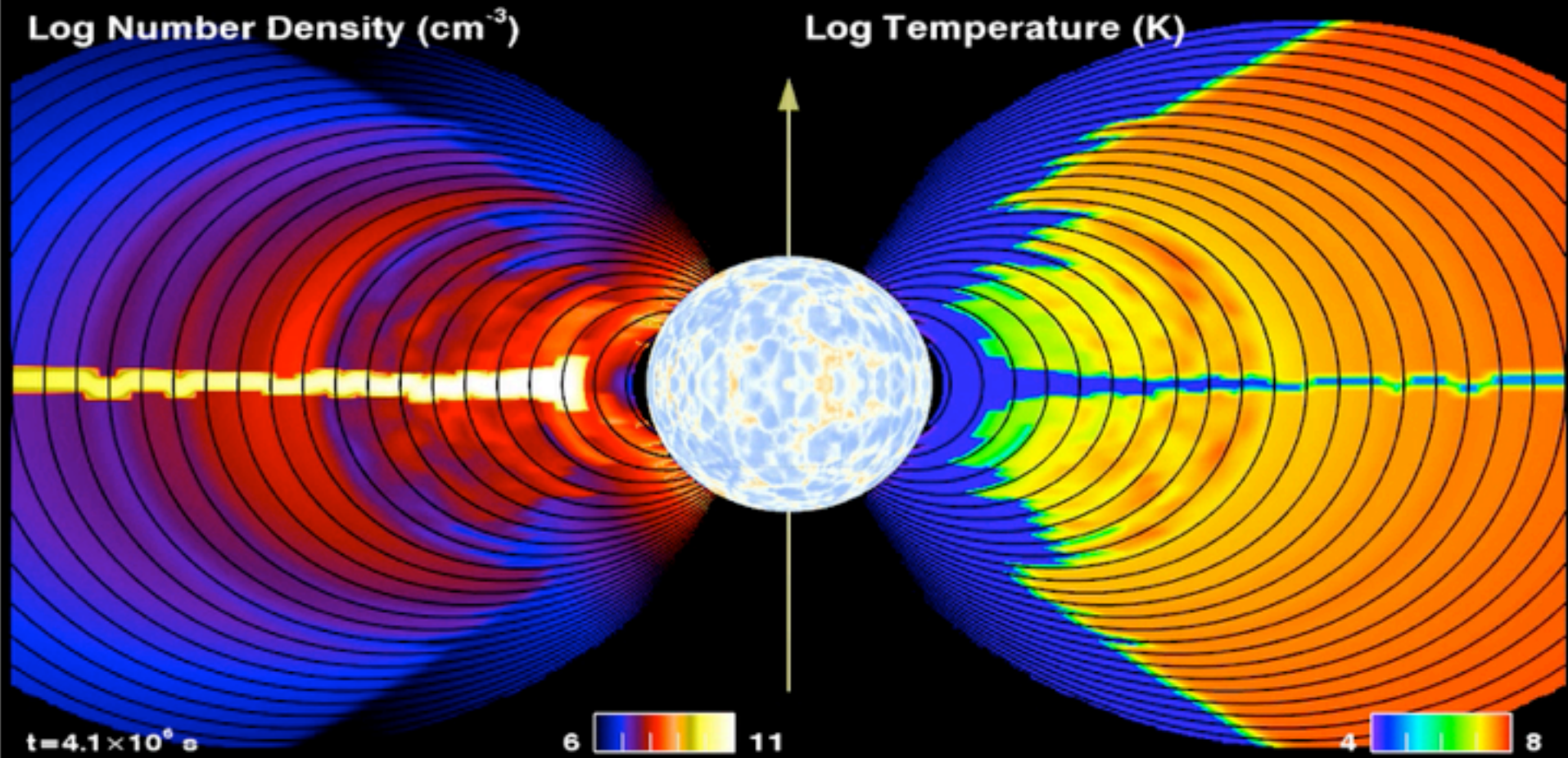
P.I.: Gregg A. Wade, Royal Military College  
50+ Co-Is, 2008-2012, CFHT Allocation: 640 hours



[http://www.physics.queensu.ca/~wade/mimes/  
MiMeS Magnetism in Massive Stars.html](http://www.physics.queensu.ca/~wade/mimes/MiMeS_Magnetism_in_Massive_Stars.html)



# Rigid Field - Hydro Model

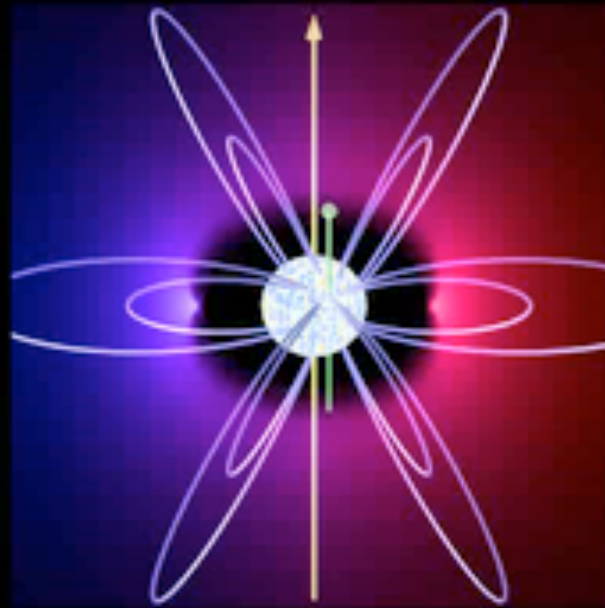




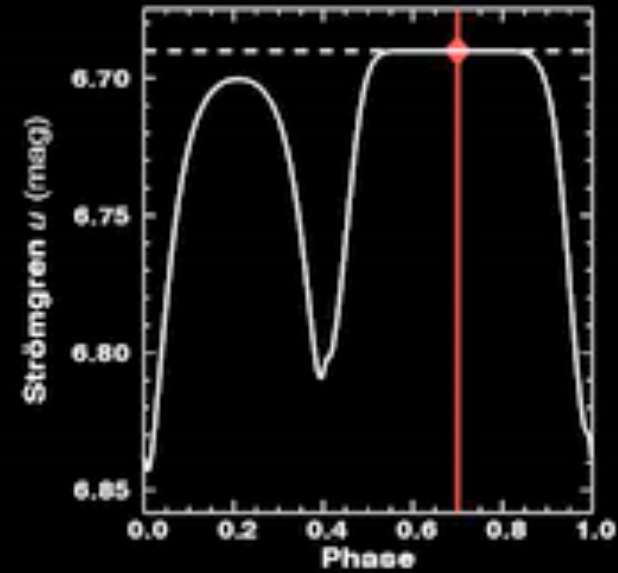
# $\sigma$ Ori E

## Rigidly Rotating Magnetosphere

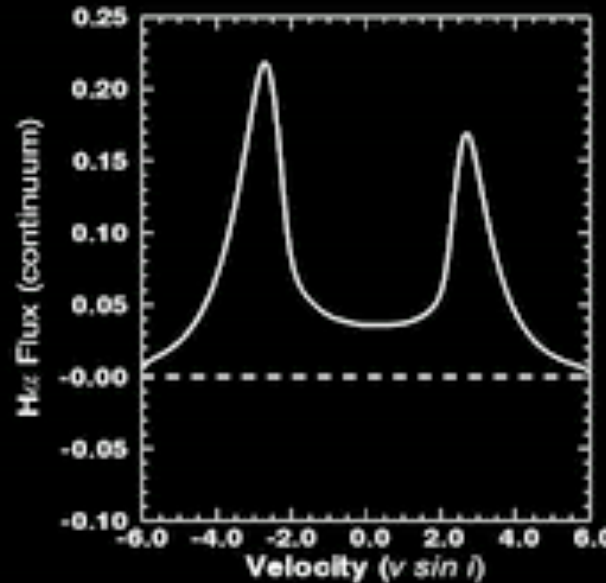
EM +B-field



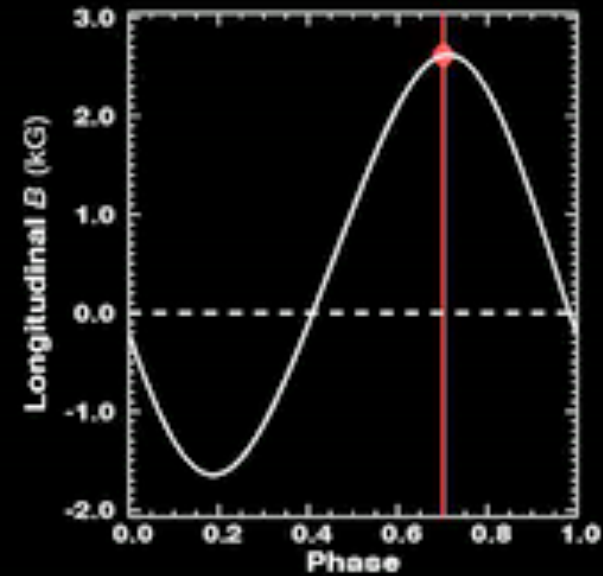
photometry



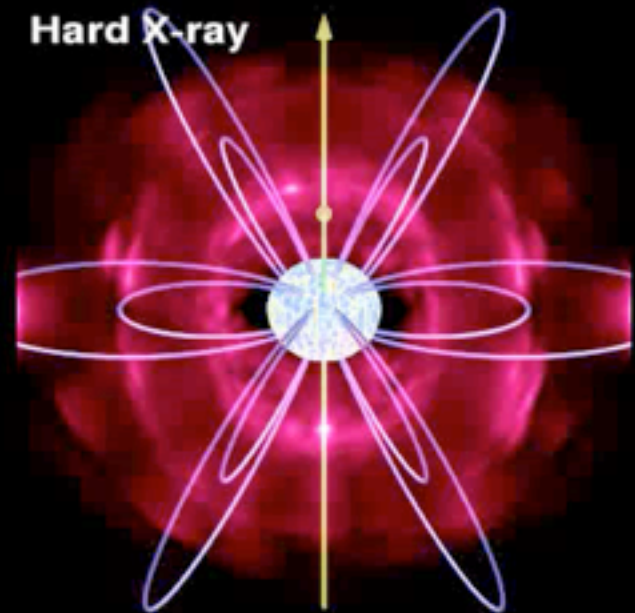
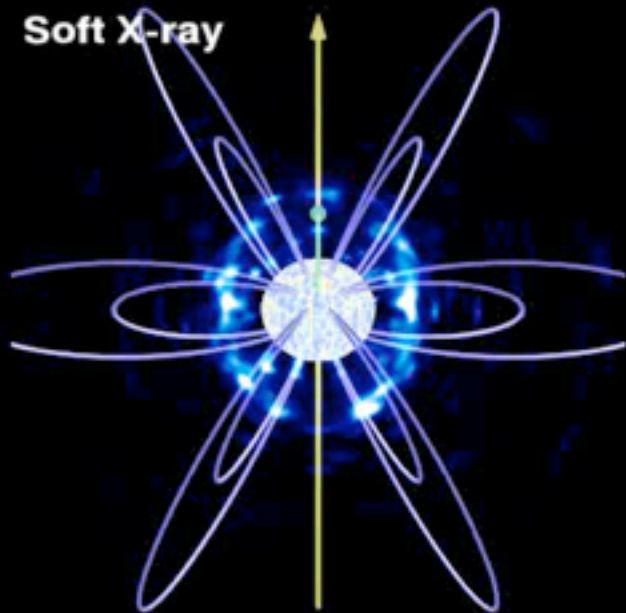
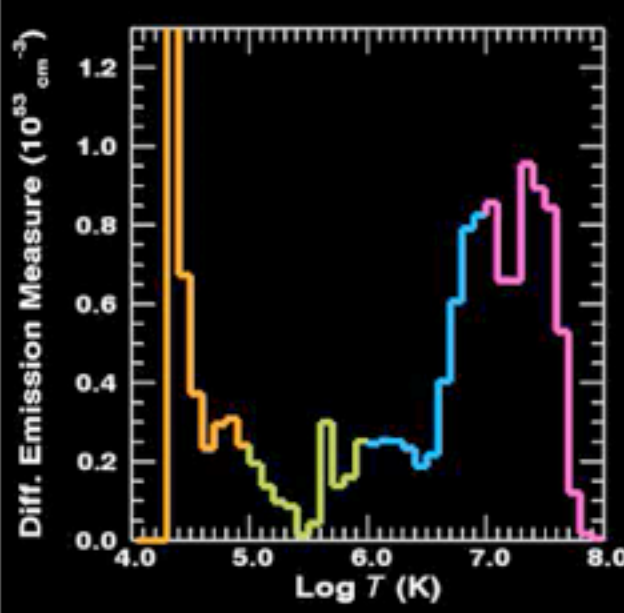
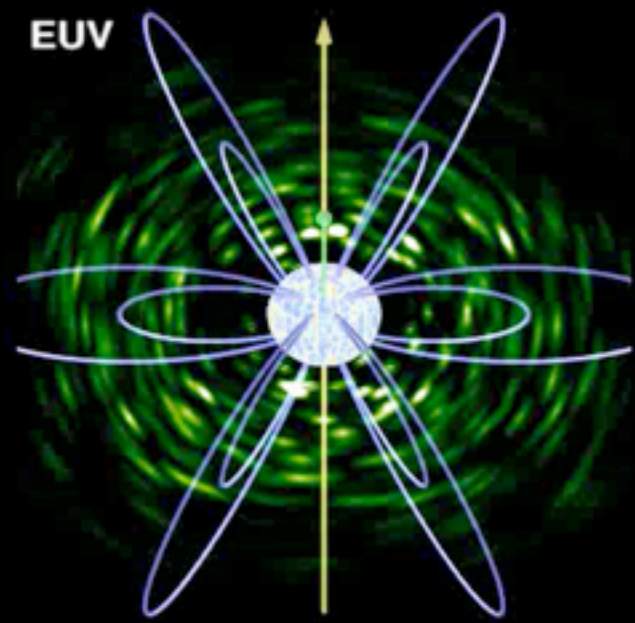
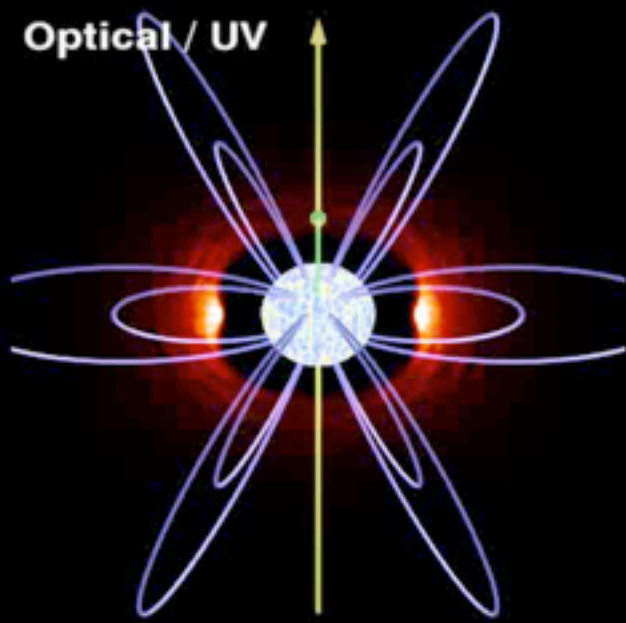
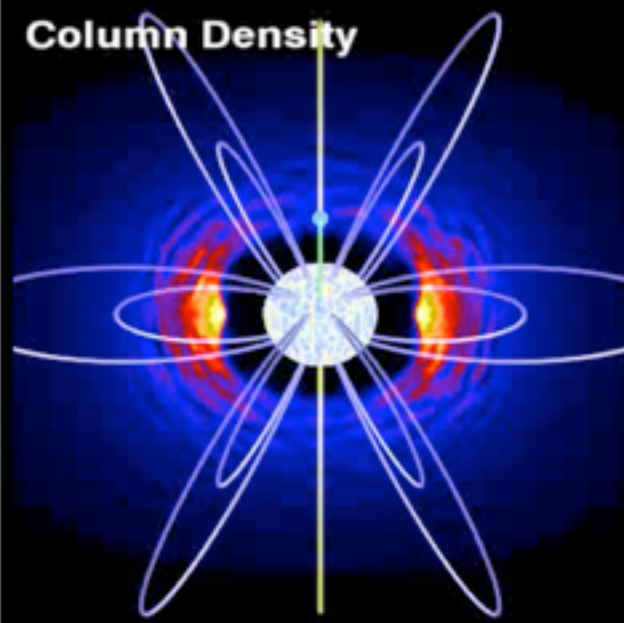
$B_* \sim 10^4$  G  
 $\Rightarrow \eta_* \sim 10^6$  !  
tilt  $\sim 55^\circ$



H $\alpha$



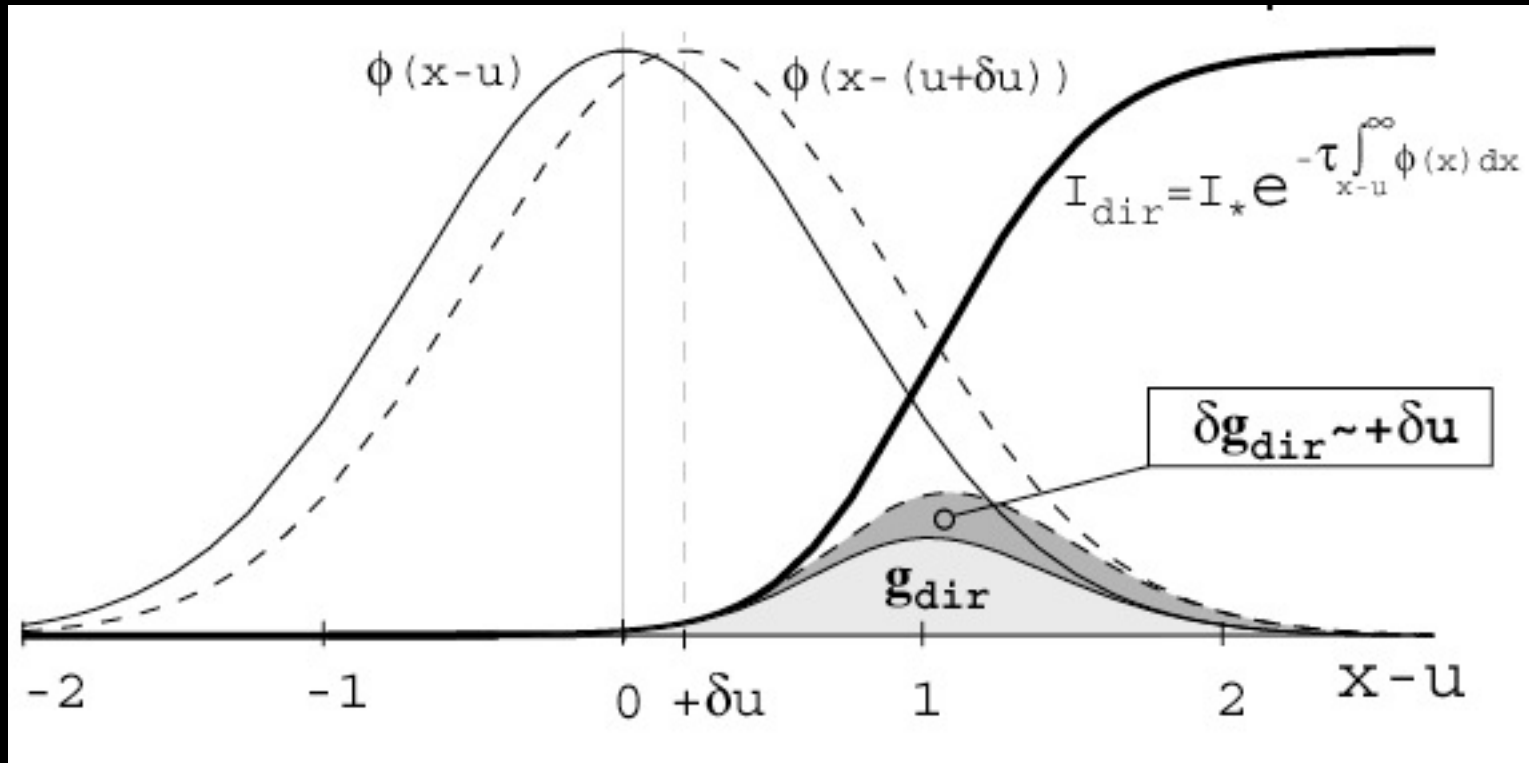
polarimetry



# Back to 1D models of non-magnetic winds & “Line-Deshadowing” Instability

- Leads to:
  - clumping
  - porosity
  - soft X-ray emission

# Line-Deshadowing Instability



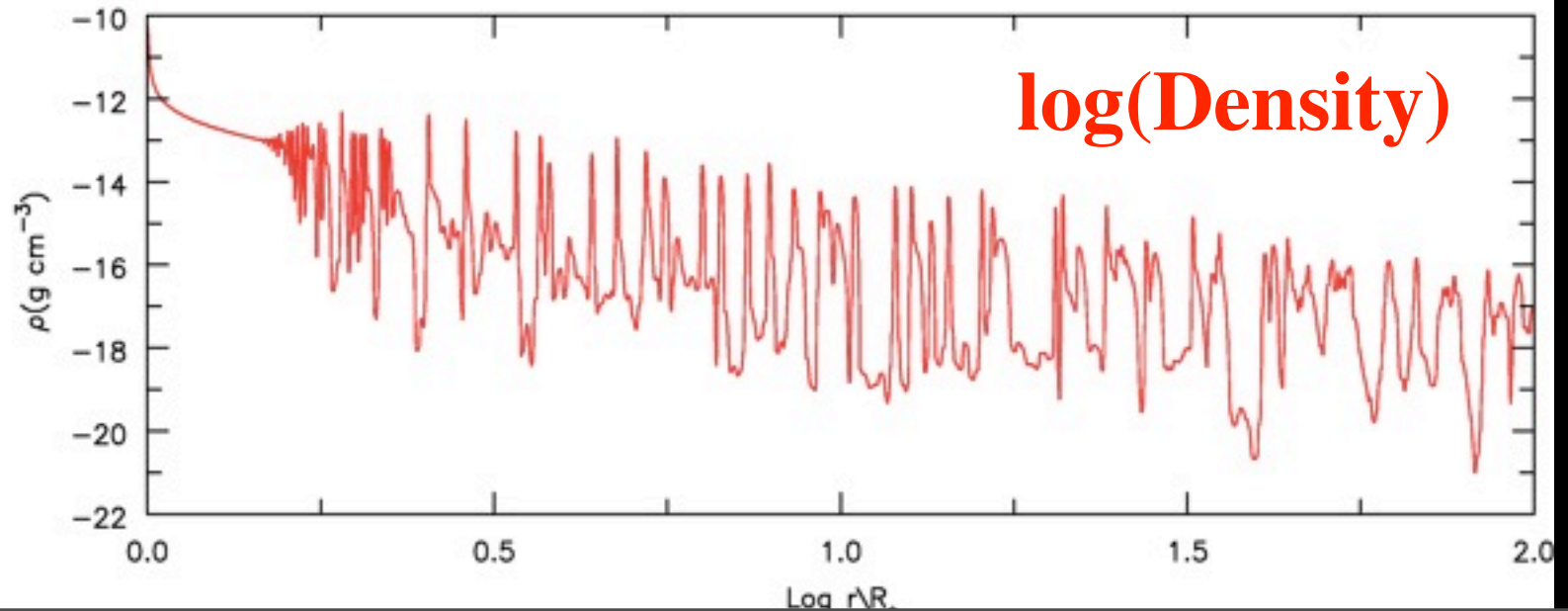
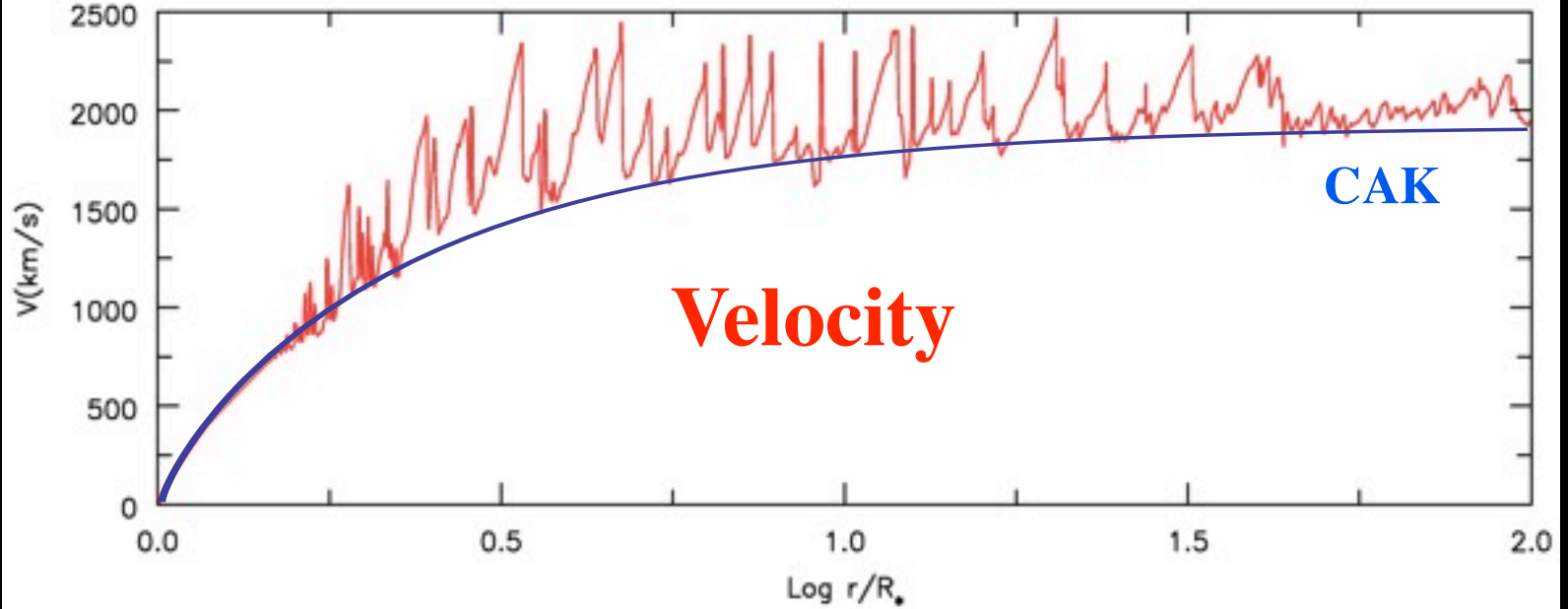
for  $\lambda < L_{\text{sob}}$ :

Instability growth rate

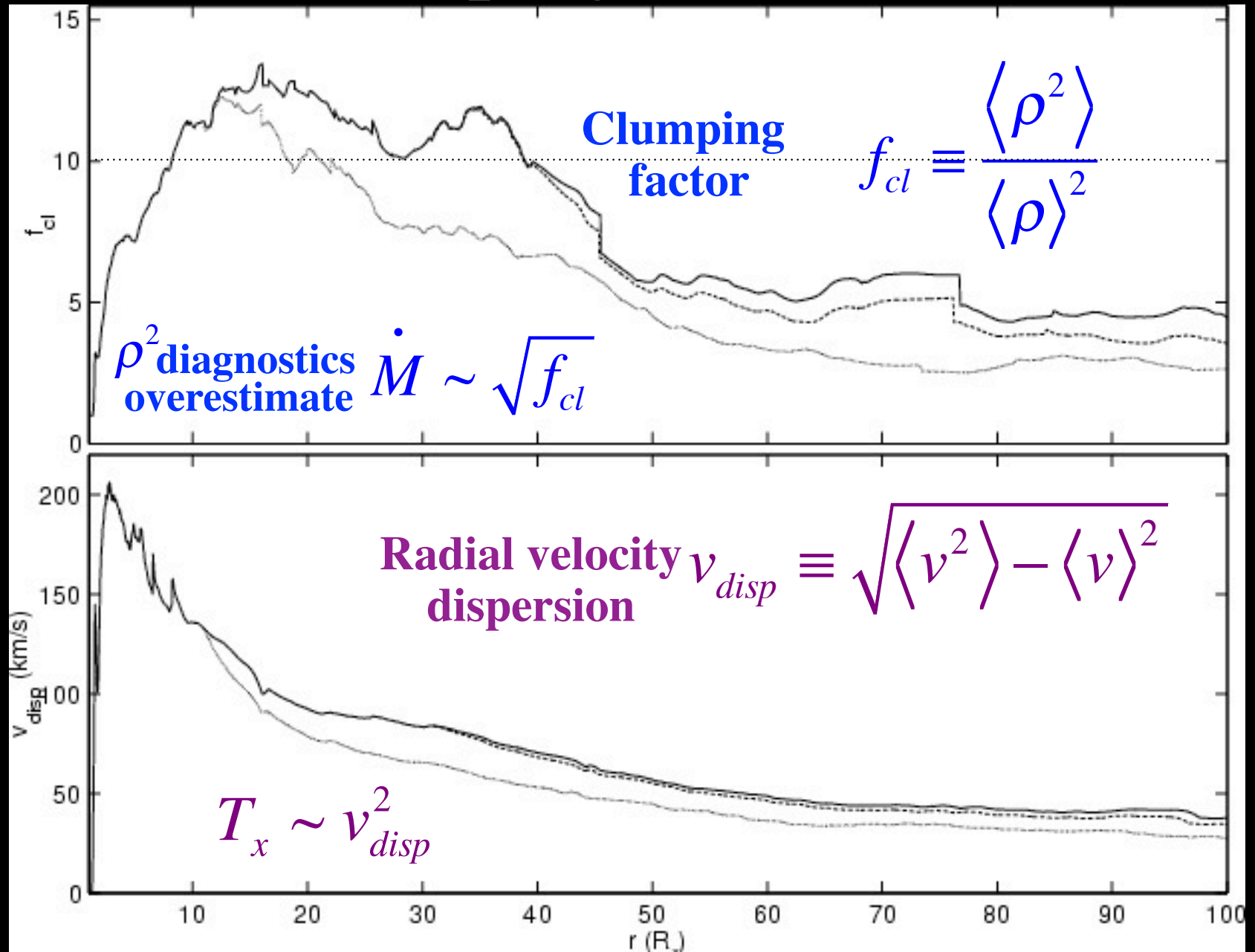
$$\Omega \sim g_0 / v_{\text{th}} \sim v v' / v_{\text{th}} \sim v / L_{\text{sob}} \sim 100 v / R$$

$\Rightarrow e^{100}$  growth!

# Time snapshot of wind structure vs. radius

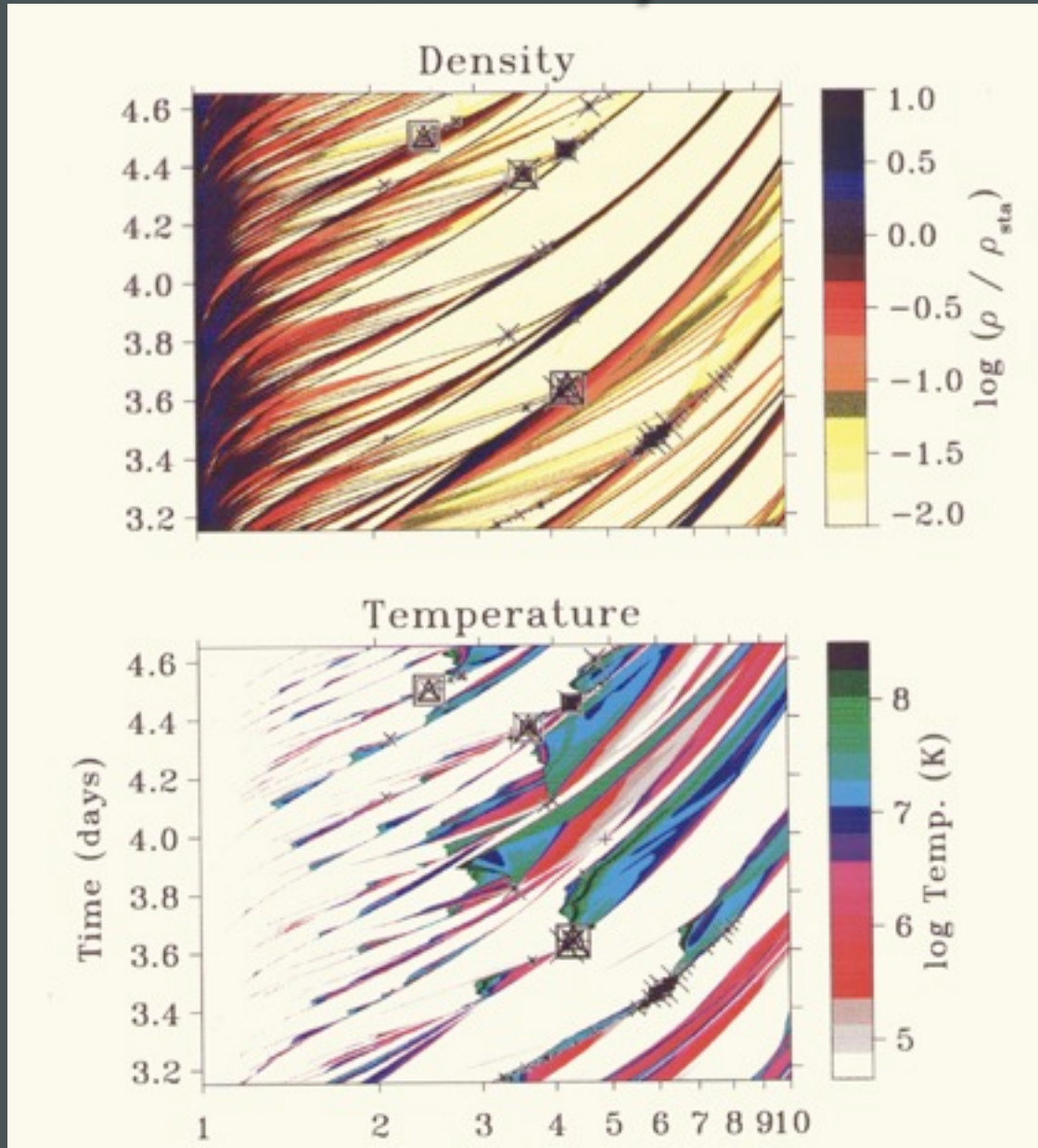


# Clumping vs. radius



# Shell-shell collisions seeded by turbulence at the base of the wind flow

Time ↑



Radius →

Feldmeier, et al. 1997

$T \sim 5-10 \text{ MK}$   
 $\Rightarrow 0.5-1 \text{ keV X-rays}$

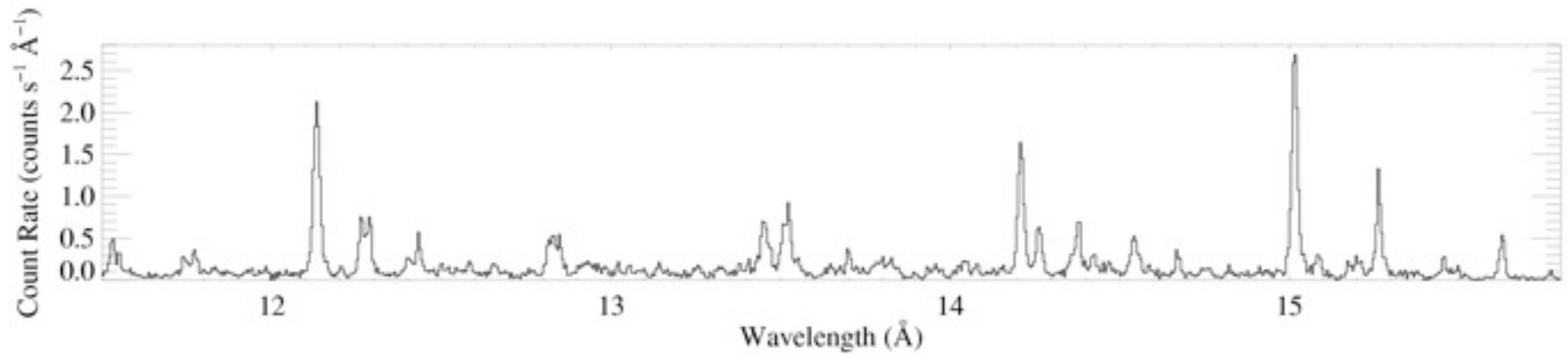
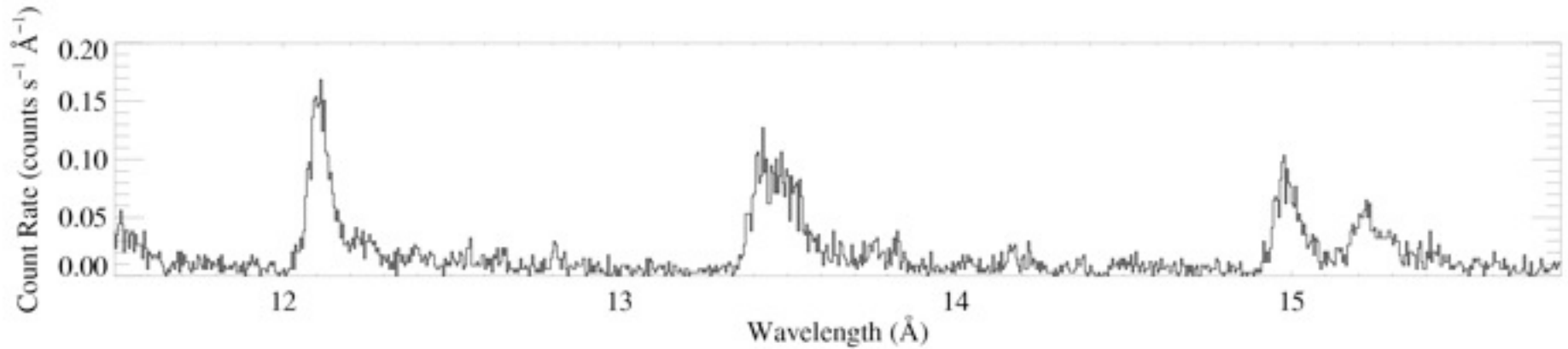
# Chandra spectra

## Hot-star $\zeta$ Pup: wind shock X-rays lines (broad)

Ne X

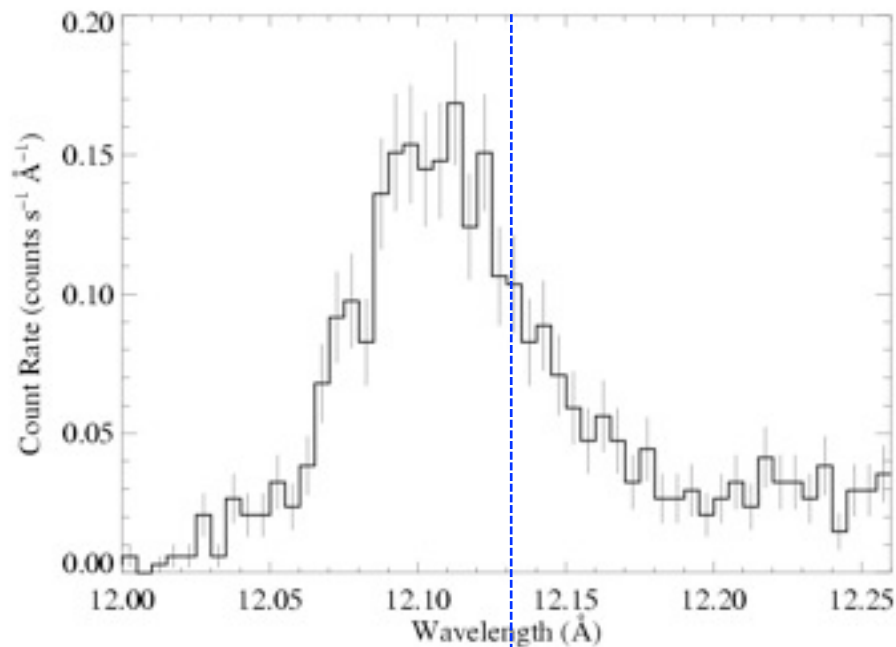
Ne IX

Fe XVII



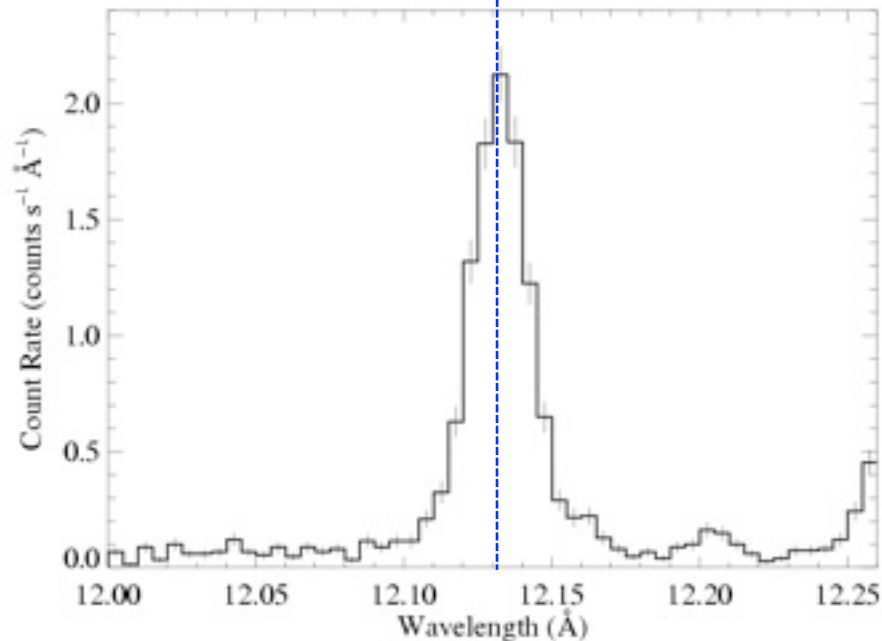
## Cool-star Capella: coronal X-rays lines (narrow)





$\zeta$  Pup  
high-mass

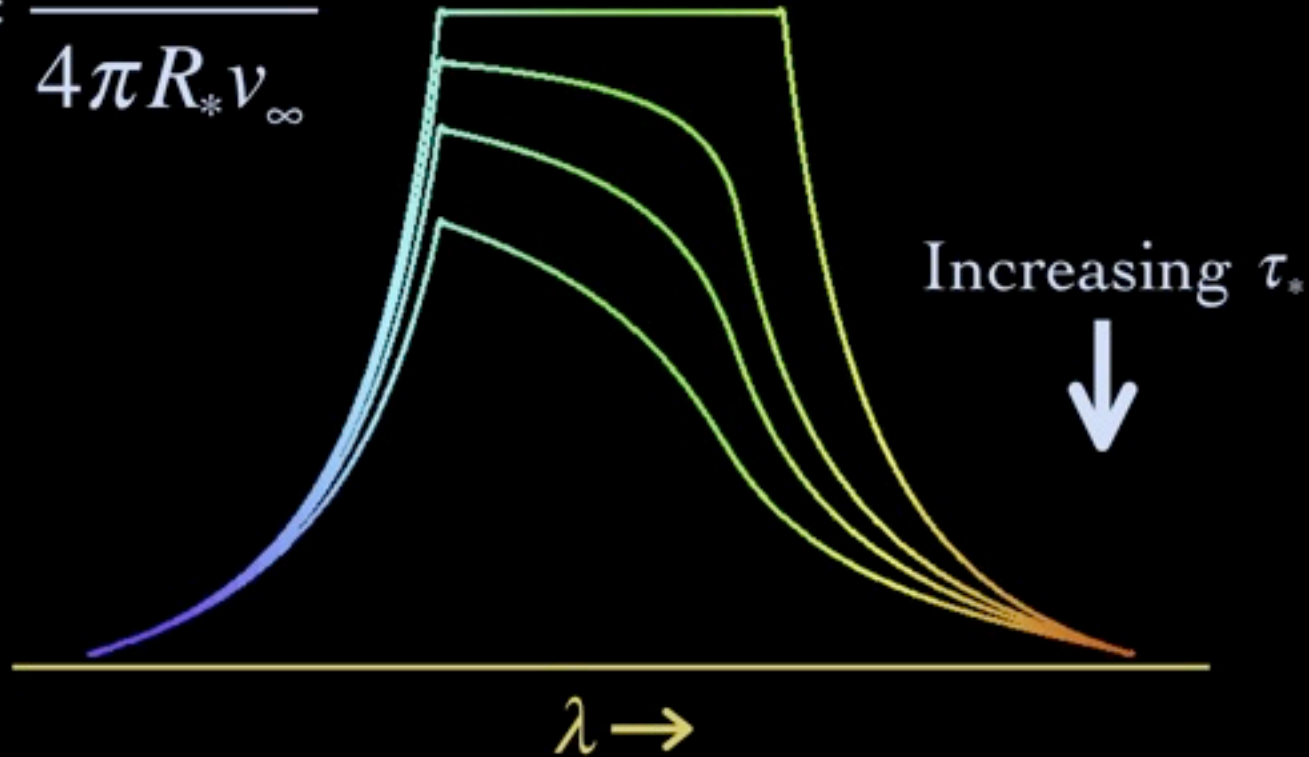
Ne X Ly  $\alpha$   
broad,  
skewed,  
blue shifted



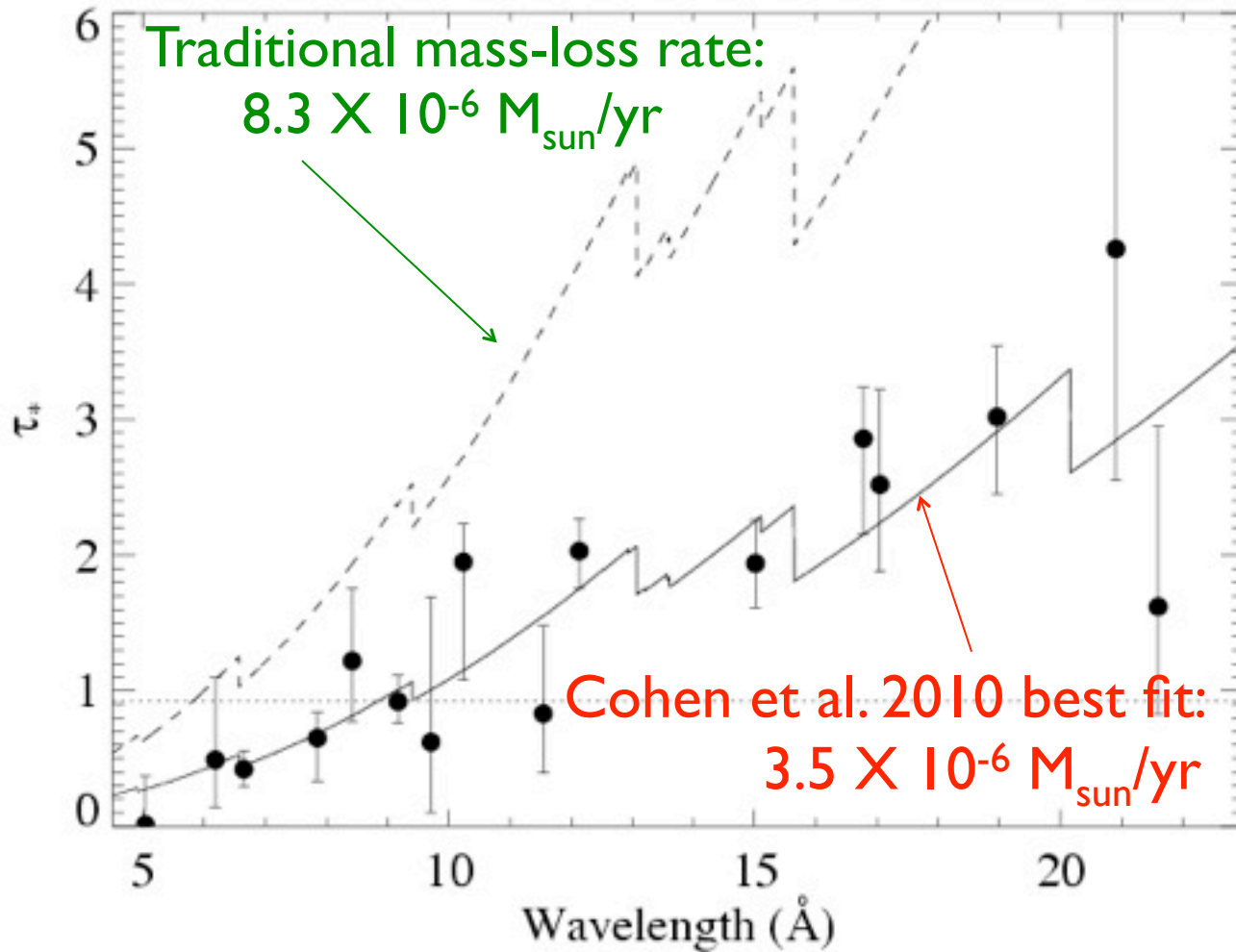
Capella  
low mass  
unresolved

# Wind Profile Model

$$\tau_* = \frac{\kappa \dot{M}}{4\pi R_* v_\infty}$$

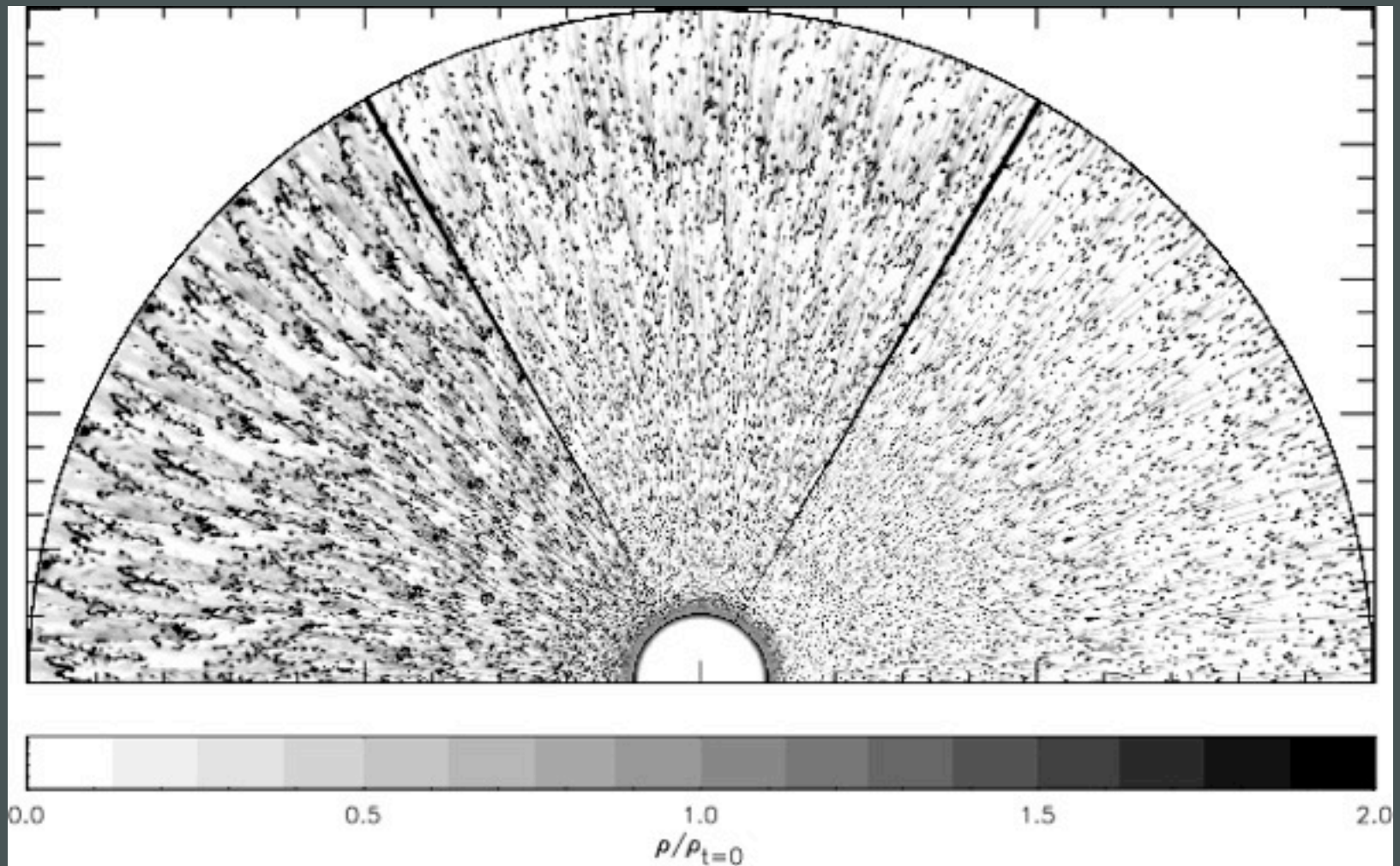


# Inferring ZPup $\dot{M}_{\text{dot}}$ from X-ray lines



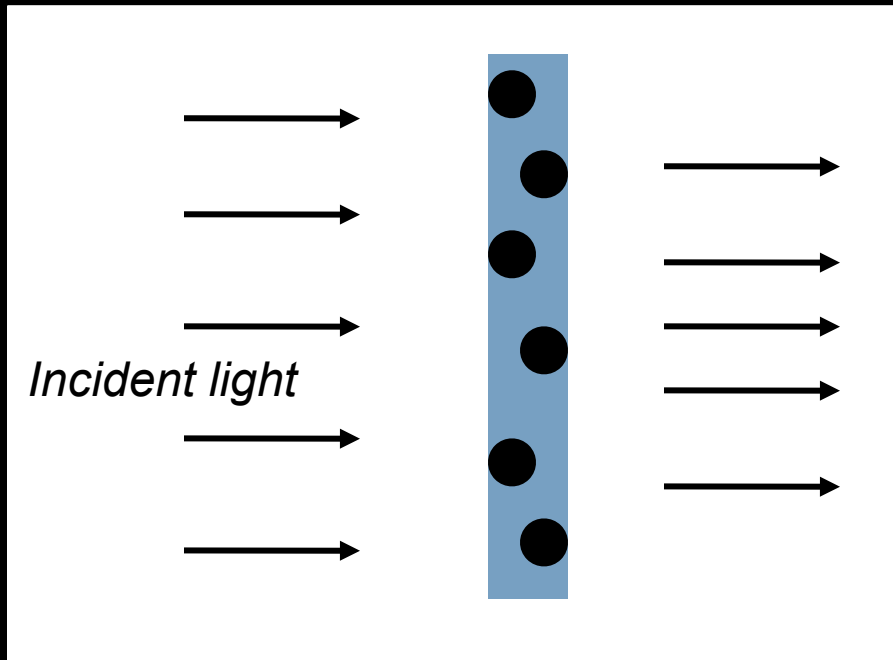
$$\dot{M} \sim \sqrt{f_{cl}} \rightarrow f_{cl} \approx 6$$

# Resulting small-scale density clumping in 2-D simulations



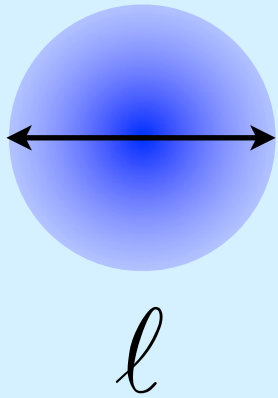
Dessart & Owocki 2003, *A&A*, 406, L1

# Porosity



- *Same amount of material*
- *More light gets through*
- *Less interaction between matter and light*

# Porous opacity from optically thick clumps



$$\sigma_{eff} \approx \ell^2 [1 - e^{-\tau_c}]$$

$$\tau_c \equiv \kappa \rho_c \ell = \kappa \rho \ell / f_{vol}$$

“porosity length”  $h$

$$K_{eff} \equiv \frac{\sigma_{eff}}{m_c} = K \frac{1 - e^{-\tau_c}}{\tau_c}$$

$$K_{eff} \approx \frac{K}{\tau_c} ; \tau_c \gg 1$$

Porosity length = mfp

$$\text{mfp} = 1/\ell^2 n_c$$

$$= L^3/\ell^2$$

$$= \ell/f_{\text{vol}}$$

$$= h = \text{porosity length}$$

clump size  $\ell = 0.05r$

$\ell = 0.1r$

$\ell = 0.2r$

Porous envelopes

$h = 0.5r$

Porosity

length

$h \equiv \ell / f_{\text{vol}}$   $h = r$

vol. fill factor

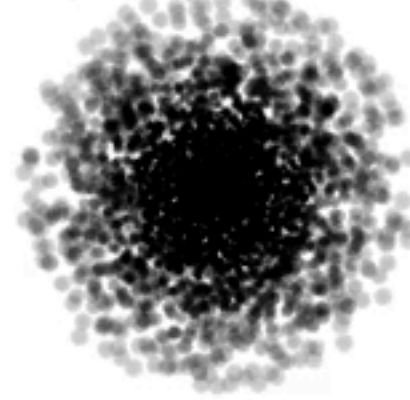
$$f_{\text{vol}} \equiv (\ell / L)^3$$

$$= 1 / f_{\text{cl}} \quad h = 2r$$

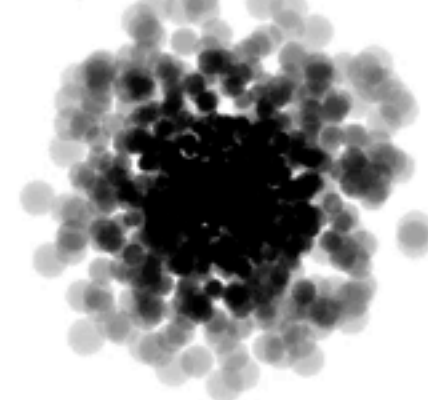
$h' = .5, \ell = .05$



$h' = .5, \ell = .1$



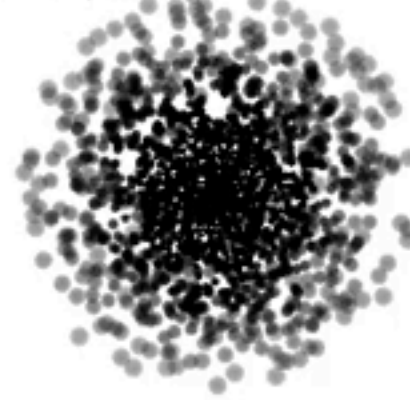
$h' = .5, \ell = .2$



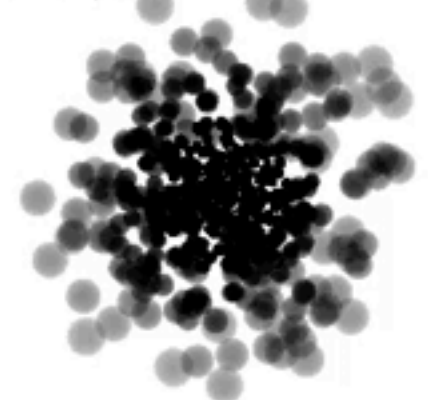
$h' = 1.0, \ell = .05$



$h' = 1.0, \ell = .1$



$h' = 1.0, \ell = .2$



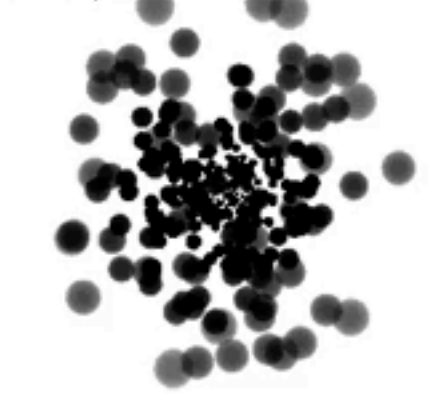
$h' = 2.0, \ell = .05$



$h' = 2.0, \ell = .1$

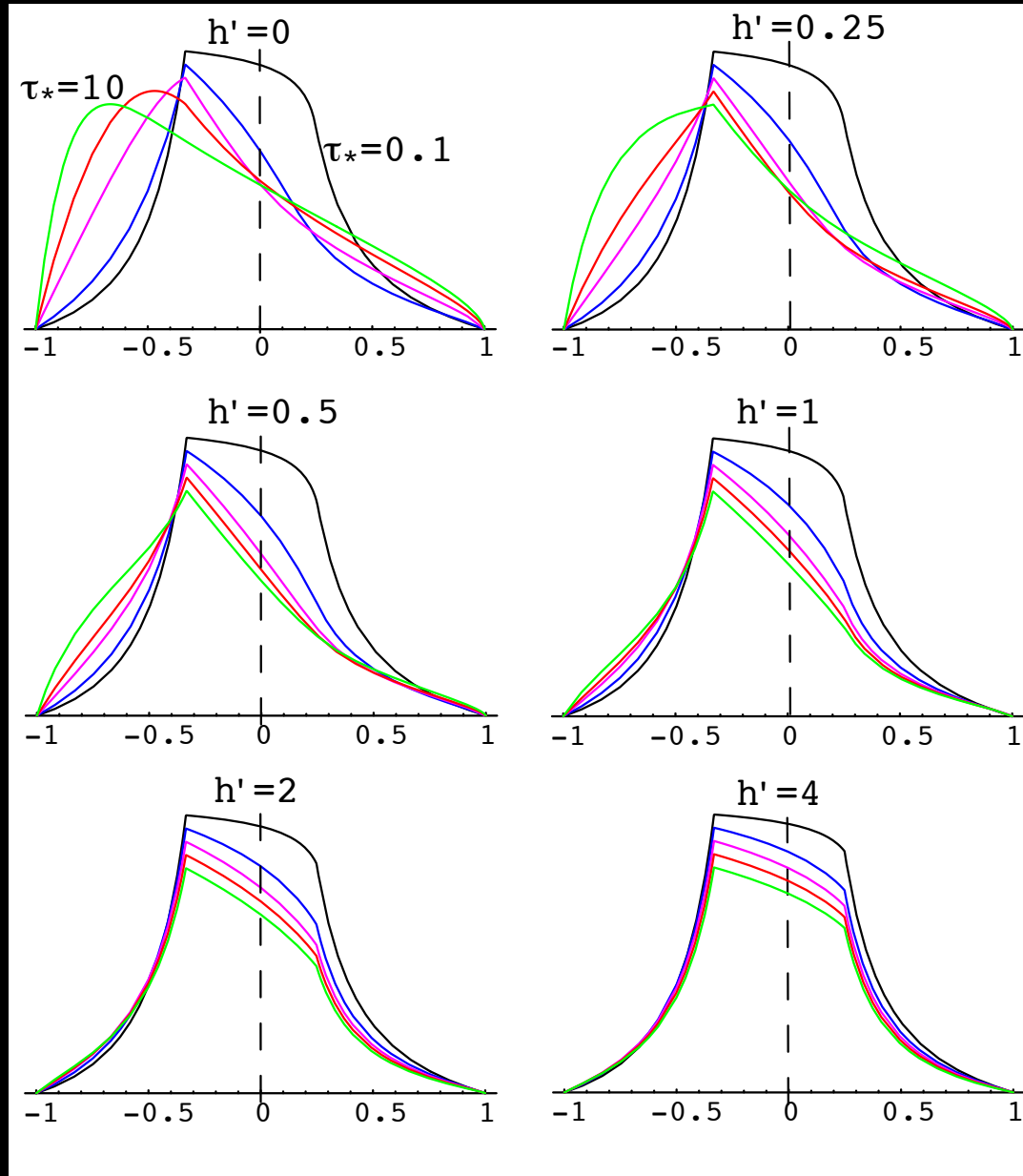


$h' = 2.0, \ell = .2$

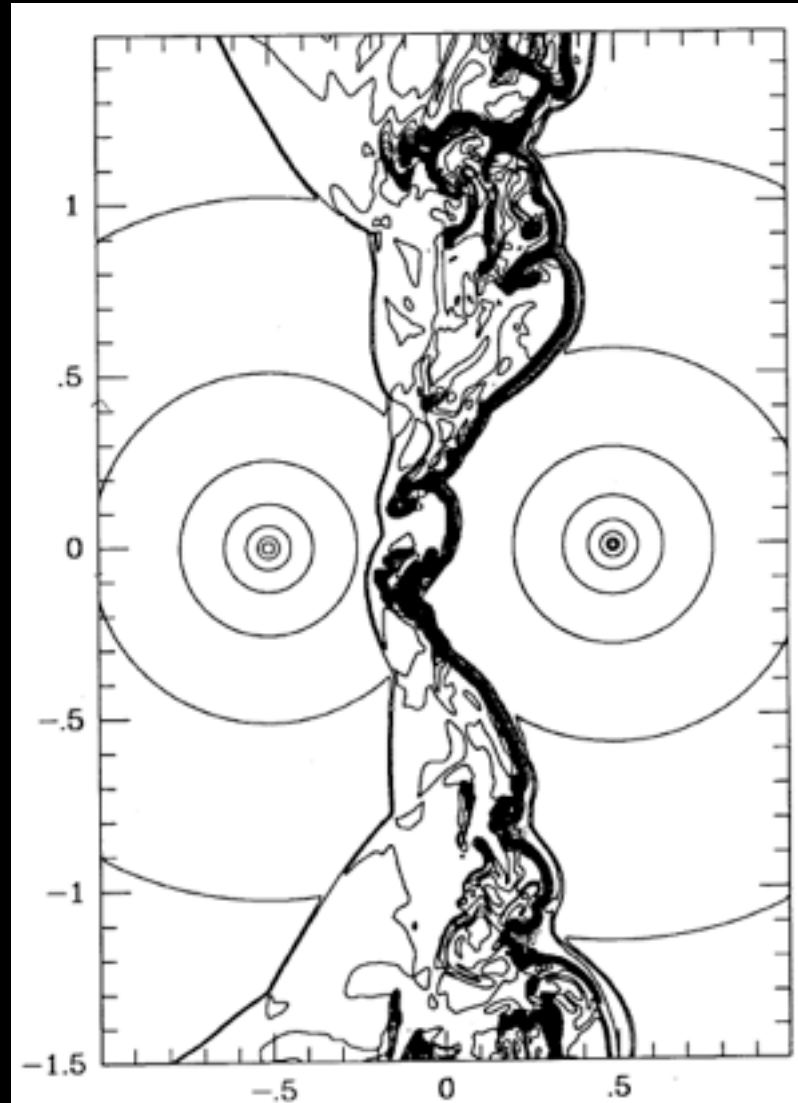




# Porosity's effect on X-ray line profiles

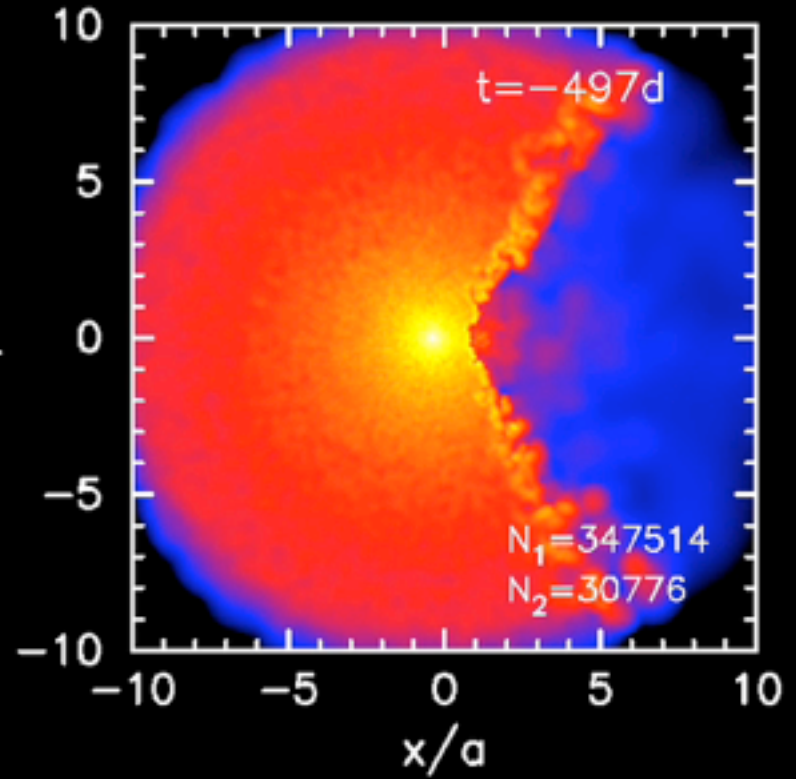
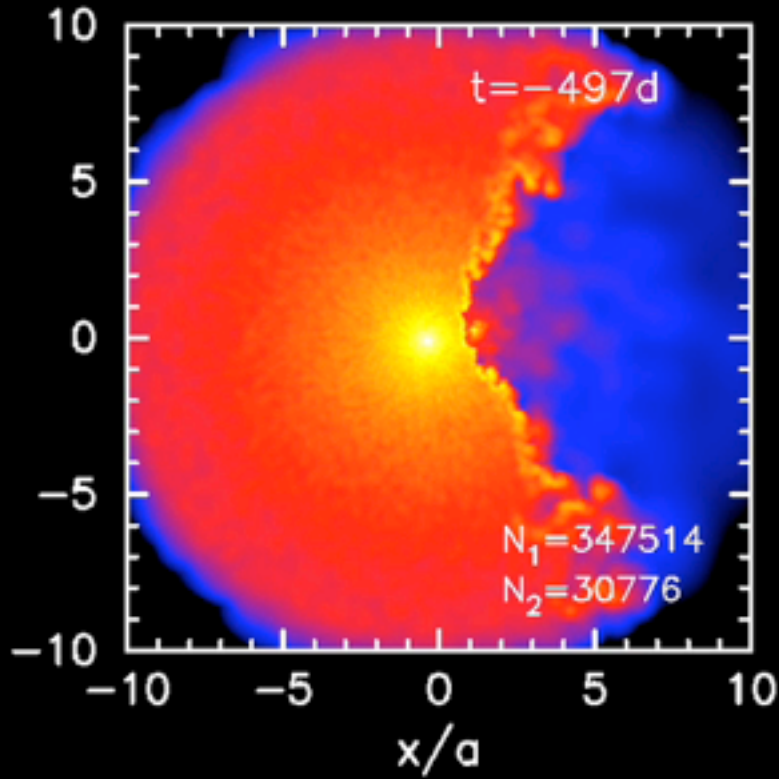


# X-rays from Colliding Wind Binaries



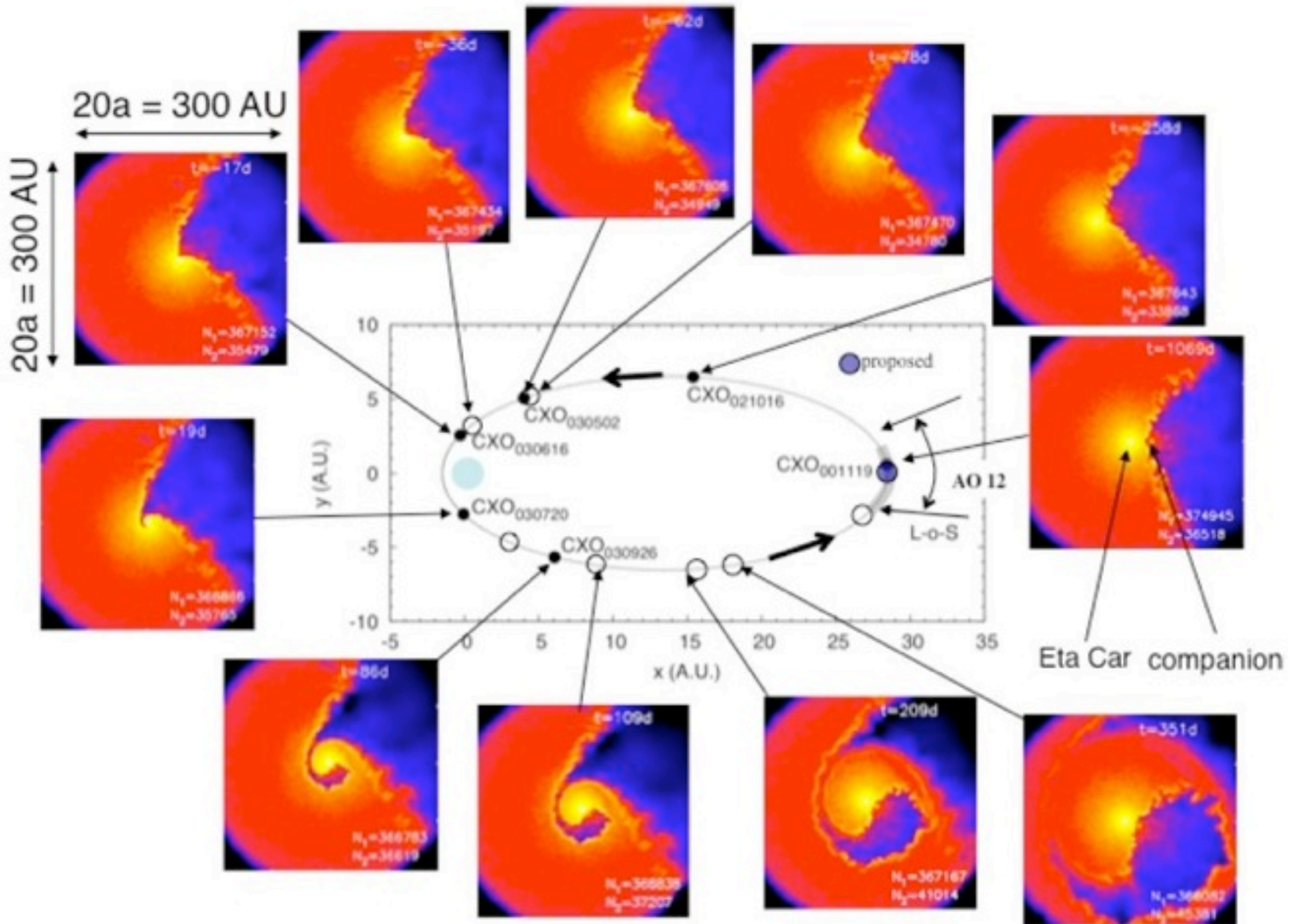
Stevens  
et al. 1992

# 3D SPH sim of CWB in eta Car

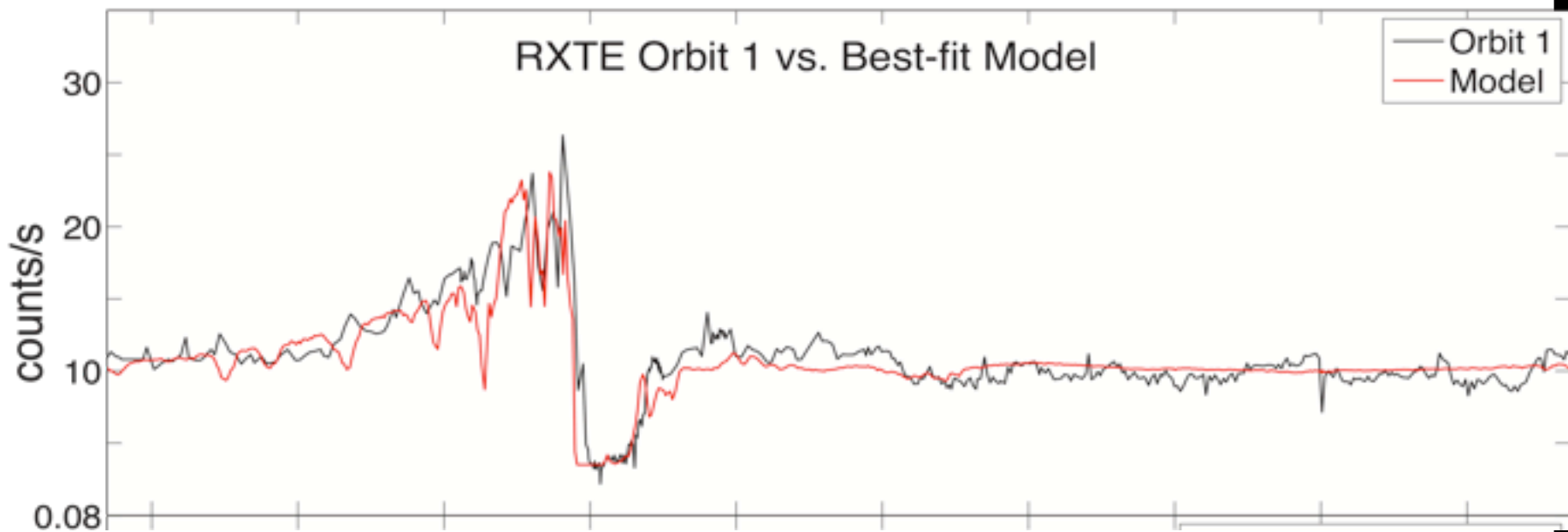
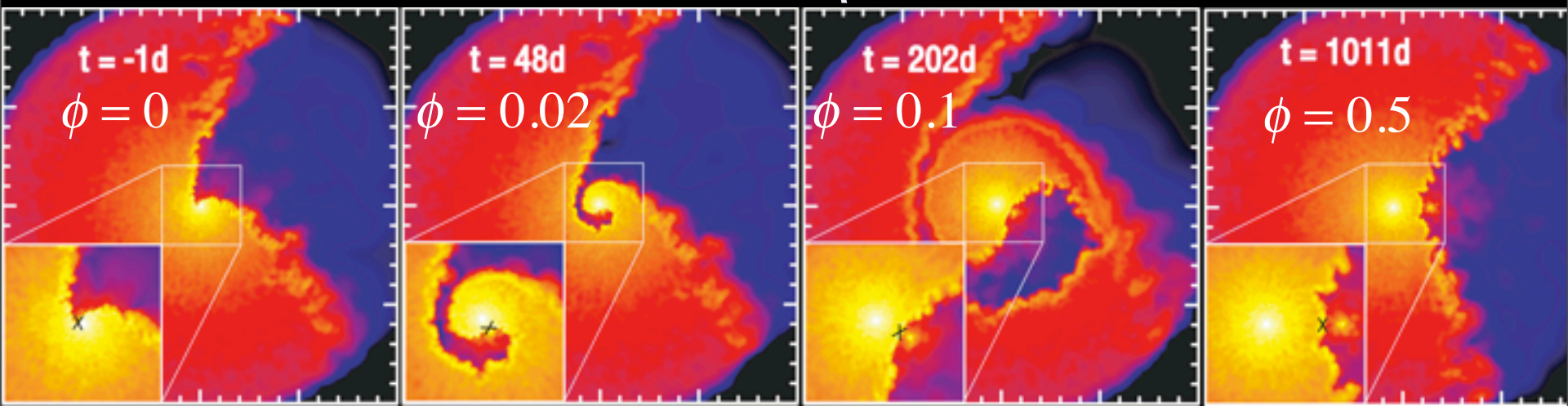


[http://www.bartol.udel.edu/~owocki/xfr/eta\\_car\\_r10-full.mov](http://www.bartol.udel.edu/~owocki/xfr/eta_car_r10-full.mov)

# 3D SPH sim of $\eta$ Car CWB

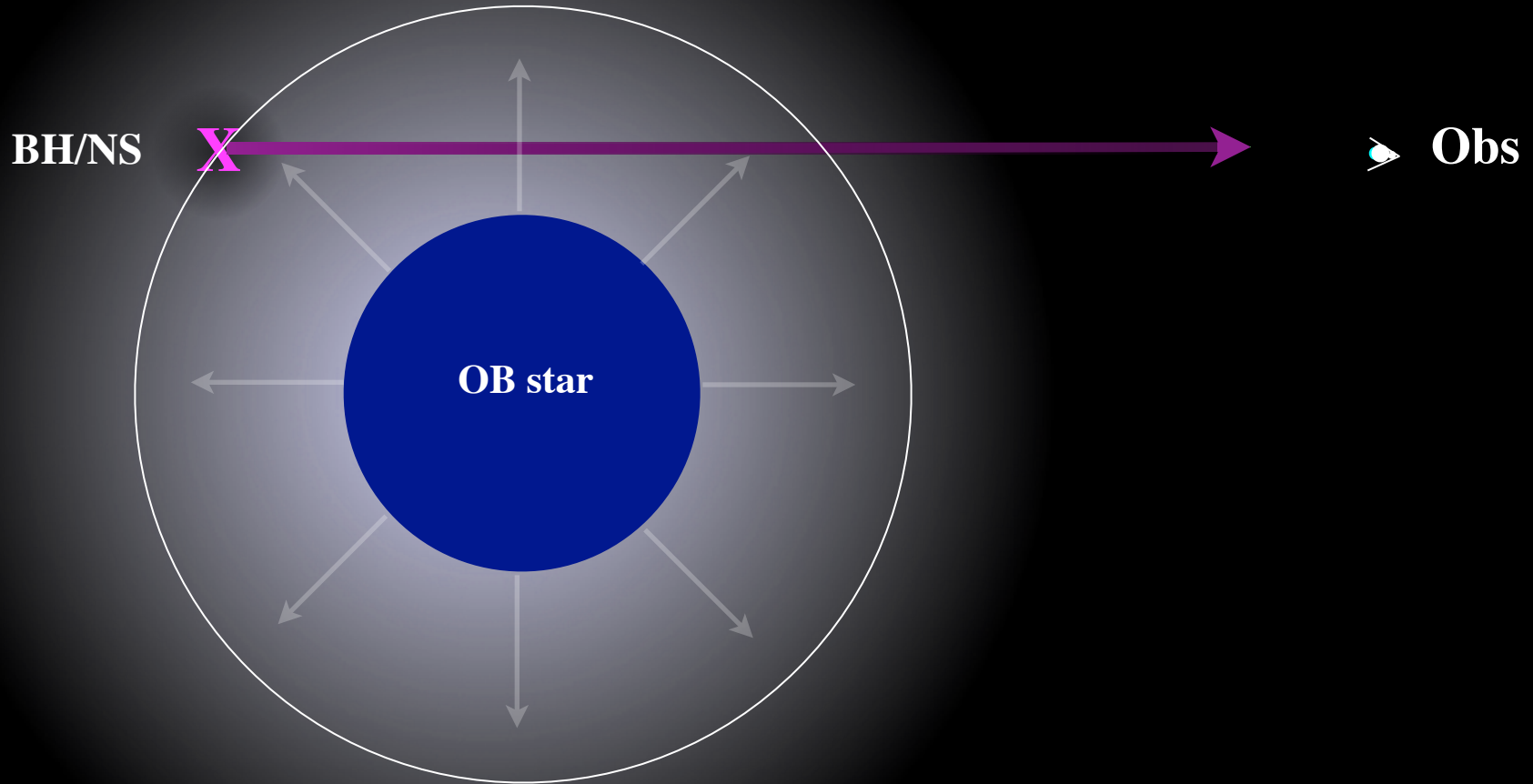


# 3D SPH sim of $\eta$ Car CWB



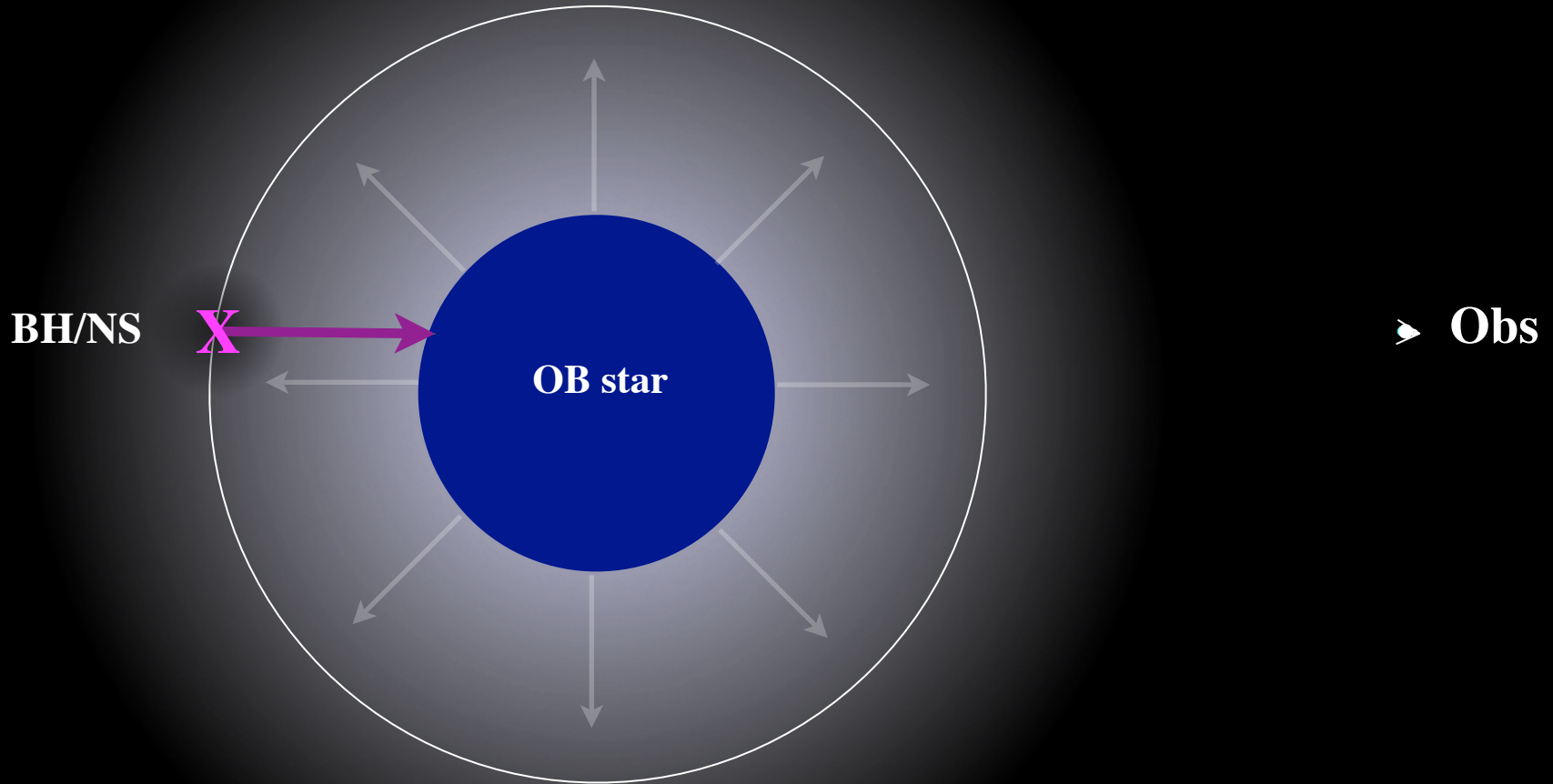
# High-Mass X-Ray Binary

**partial absorption by stellar wind**



# High-Mass X-Ray Binary

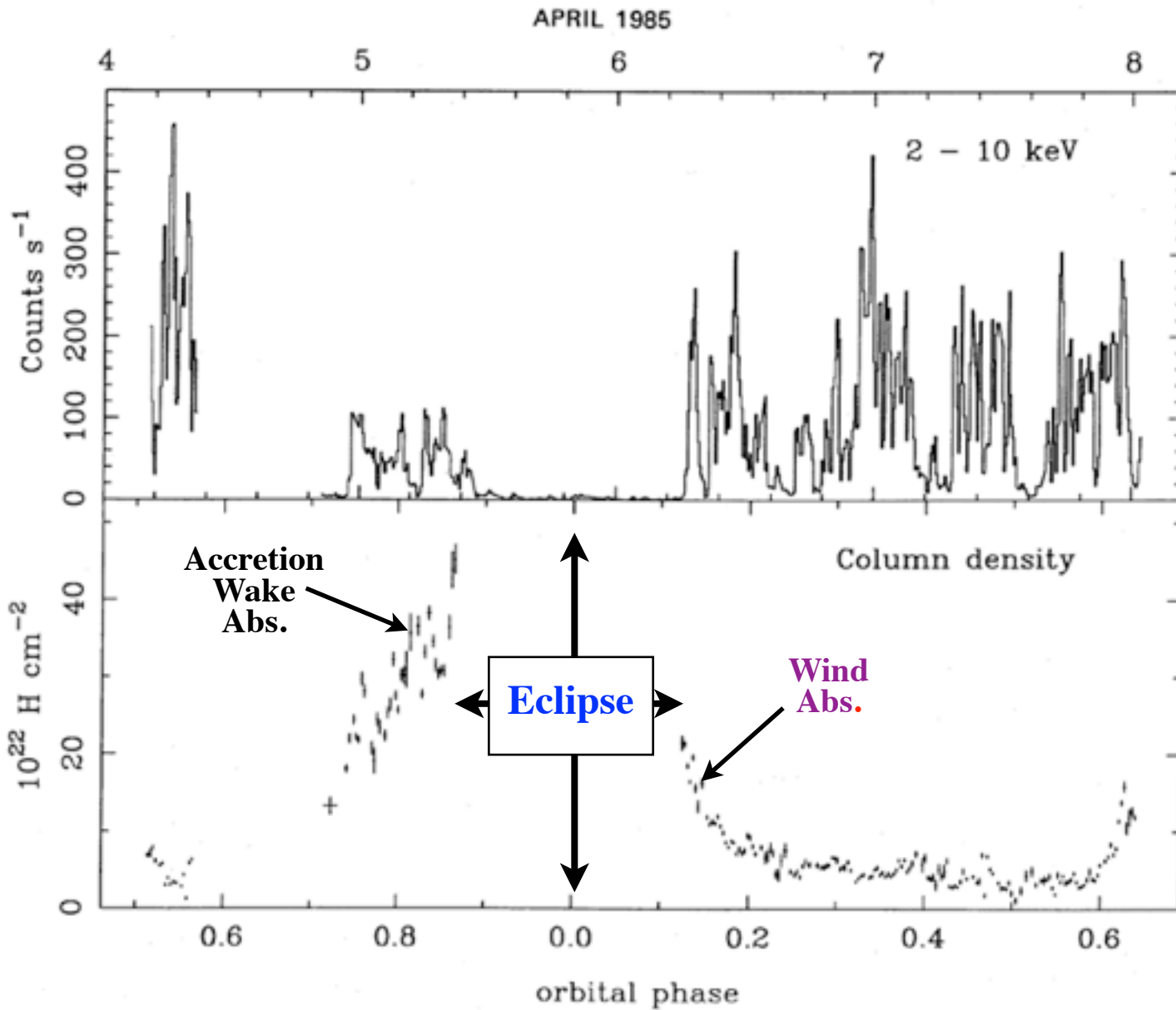
## full absorption by stellar Eclipse



# HMXRB X-ray light curve

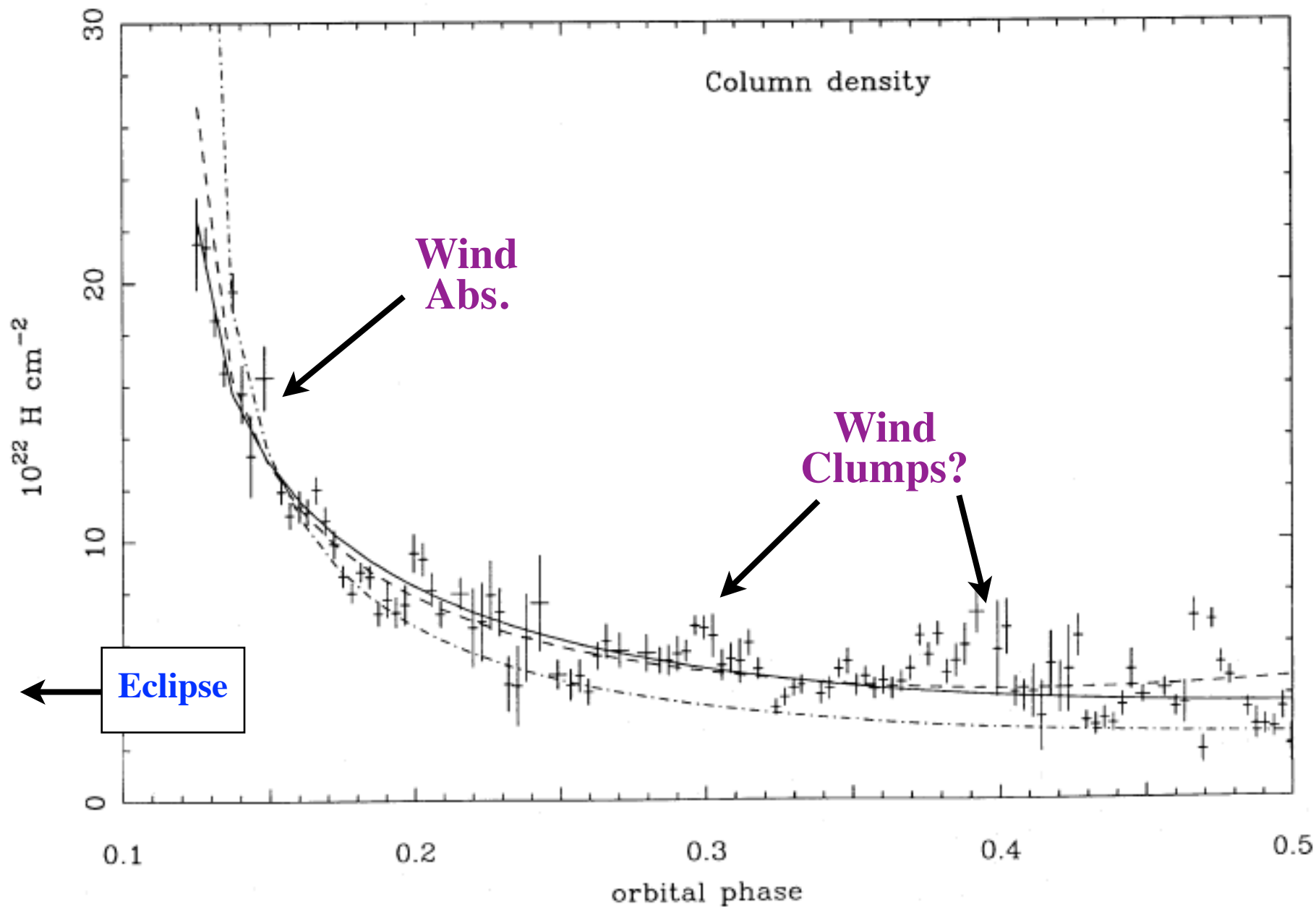
4U 1700-37  
HD 153919

Haberl et al. 1989





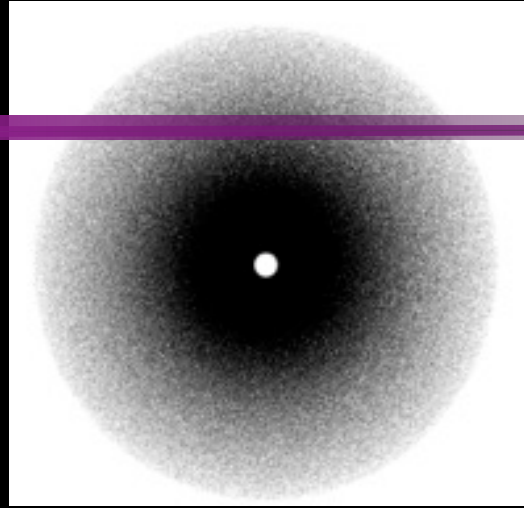
# HMXRB light curve fluctuations



# X-ray absorption in smooth vs. porous wind

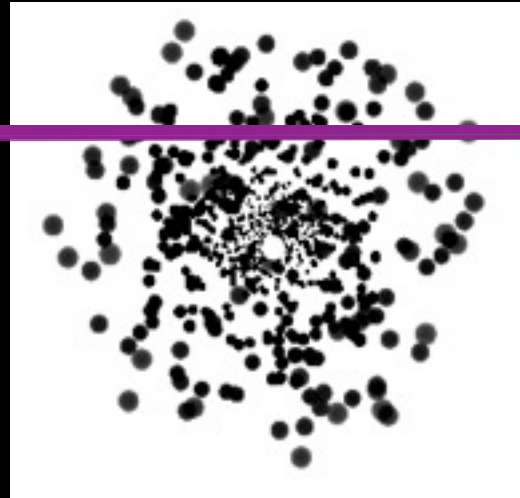
BH/NS

X

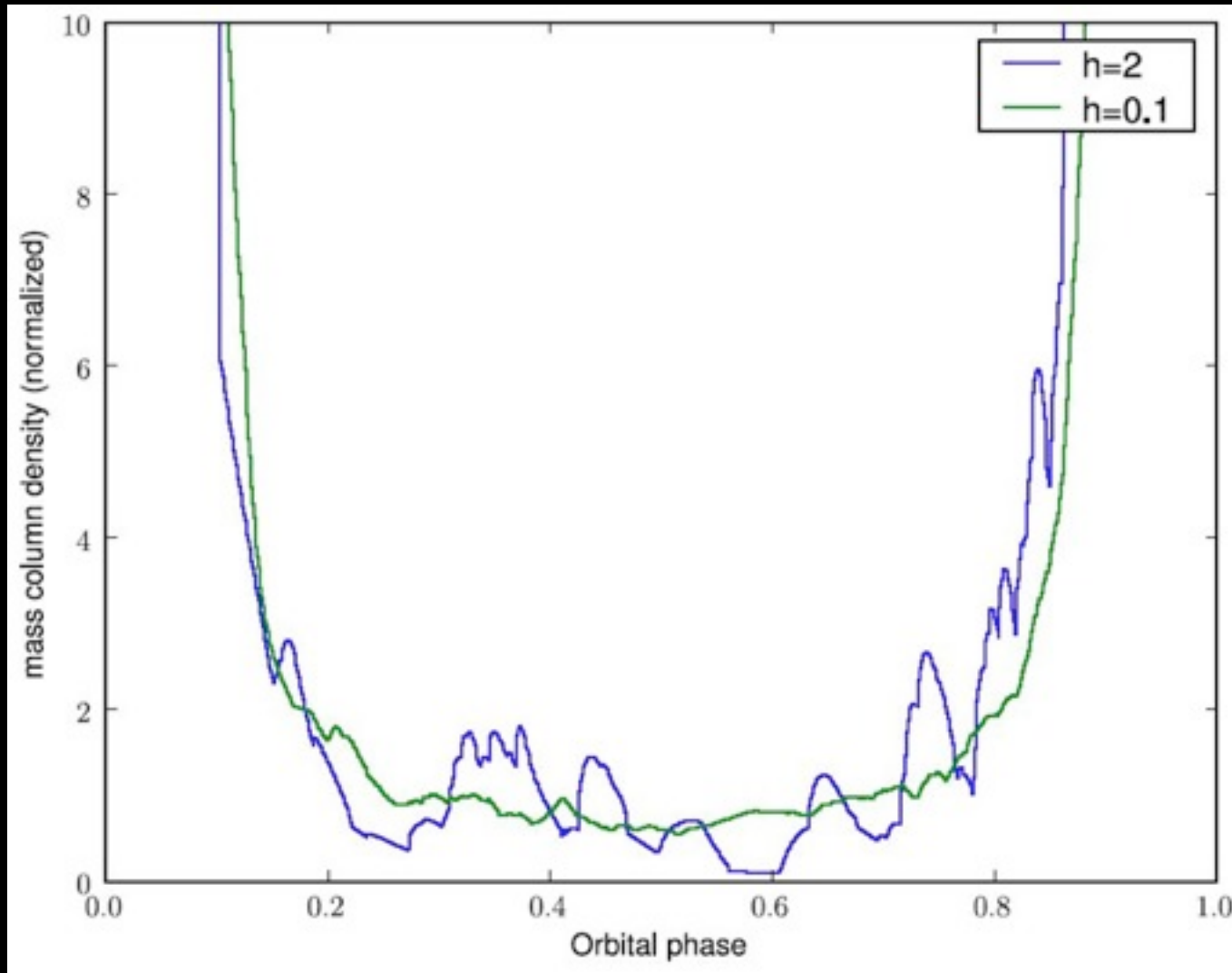


BH/NS

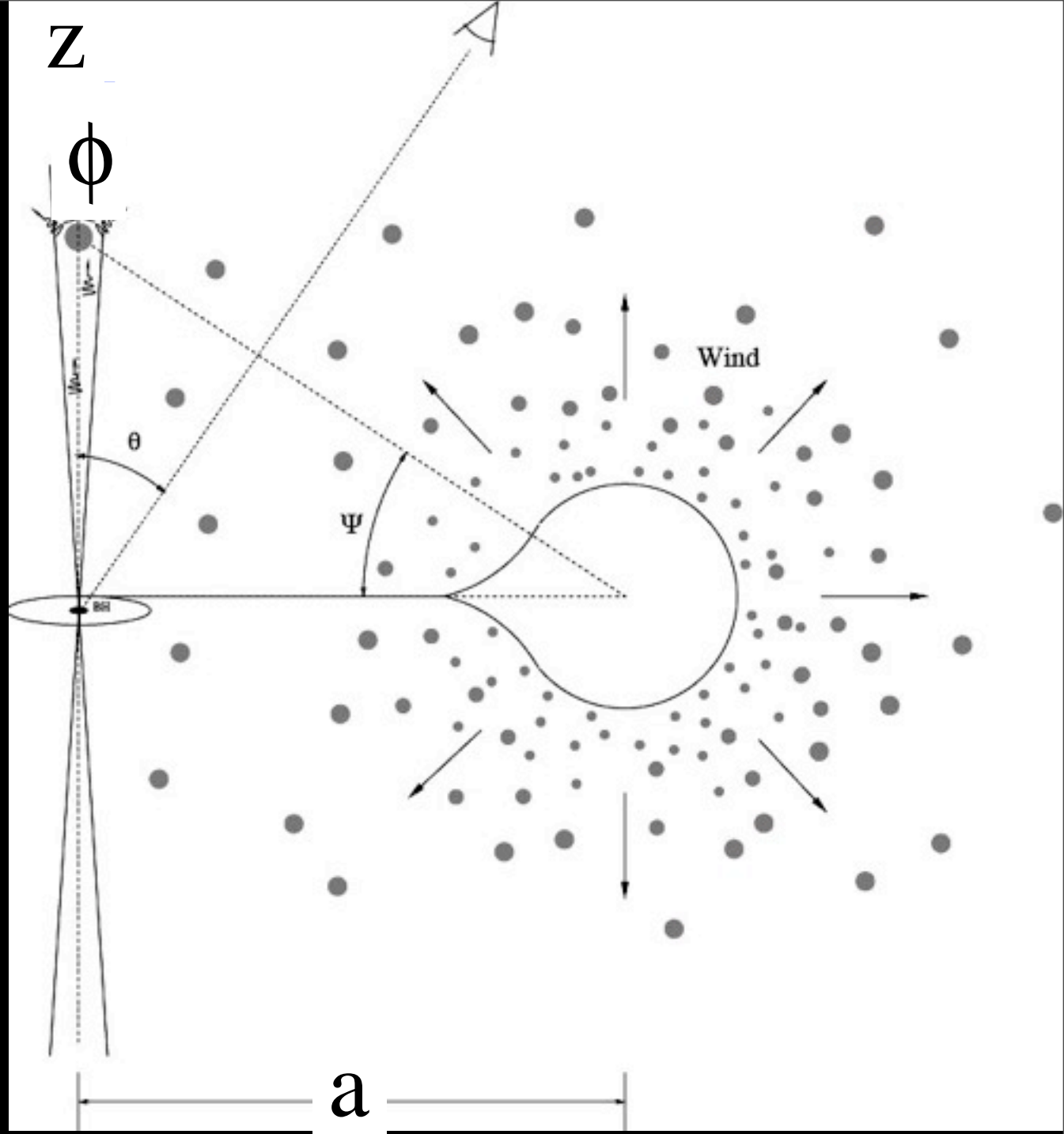
X



# Porosity Model for how wind clumps perturb light curve mass column



$\gamma$ -ray  
emission  
in  
HMXRB  
with  
clumpy  
wind



# $\gamma$ -ray fluctuation from wind clumps

$$L_\gamma = L_j \sigma \int_0^\infty n(z) dz \quad \# \text{ clumps} \quad \Delta N_c = \frac{\Delta z}{\boxed{h}} \text{ mfp}$$

$$\delta L_\gamma \sim \sqrt{N_c}$$

$$\frac{\delta L_\gamma}{L_\gamma} = \frac{\sqrt{\int_0^\infty n^2 h dz}}{\int_0^\infty n dz} = \sqrt{\frac{h}{\pi a}}$$

# Typical example

narrow jet with:  $l=h/10=0.03a$

$$\frac{\delta L_\gamma}{L_\gamma} \approx \sqrt{\frac{h}{\pi a}} \approx 0.1 = 10\%$$

# Modeling TeV gamma-rays from LS 5039: An Active OB Star at the Extreme

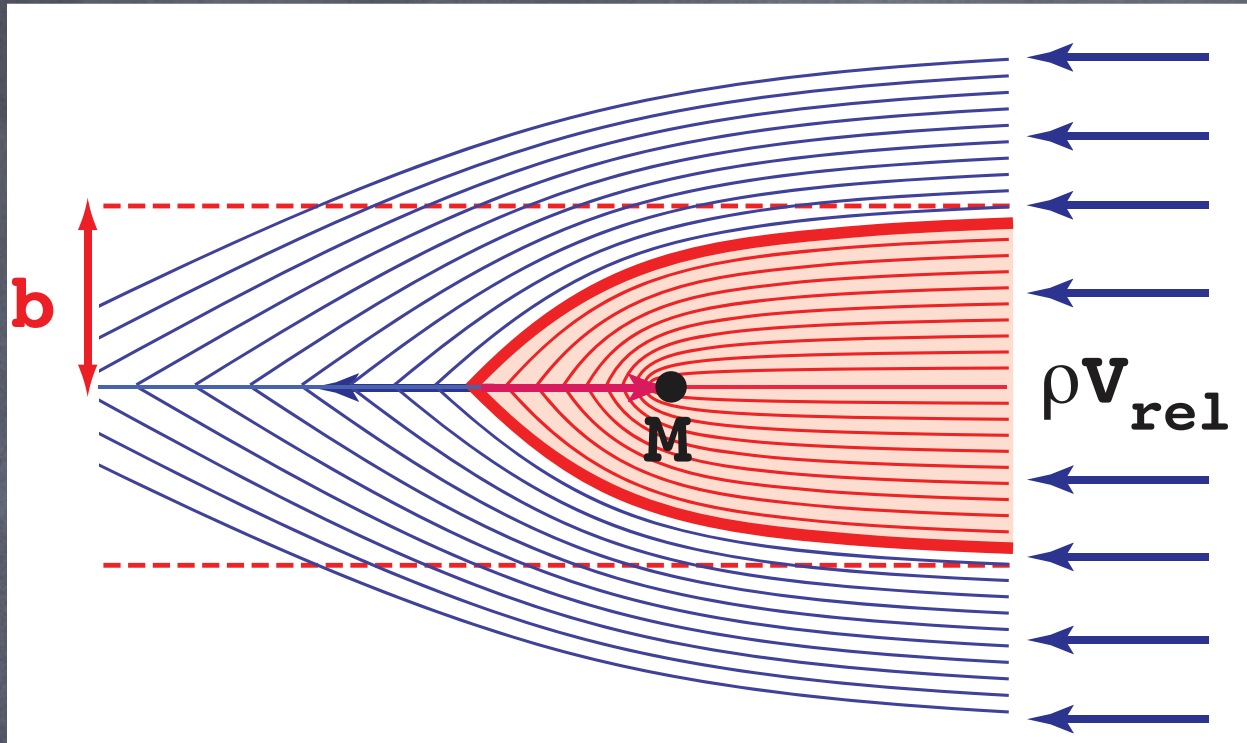
Stanley P. Owocki (BRI, Univ. Delaware)

Atsuo T. Okazaki (Hokkai-Gakuen Univ.)

Gustavo E. Romero (U. Nacional La Plata)

talk at Paris IAUS 272 July 2010

# Bondi-Hoyle-Lyttleton (BHL) Accretion



**Bondi radius**

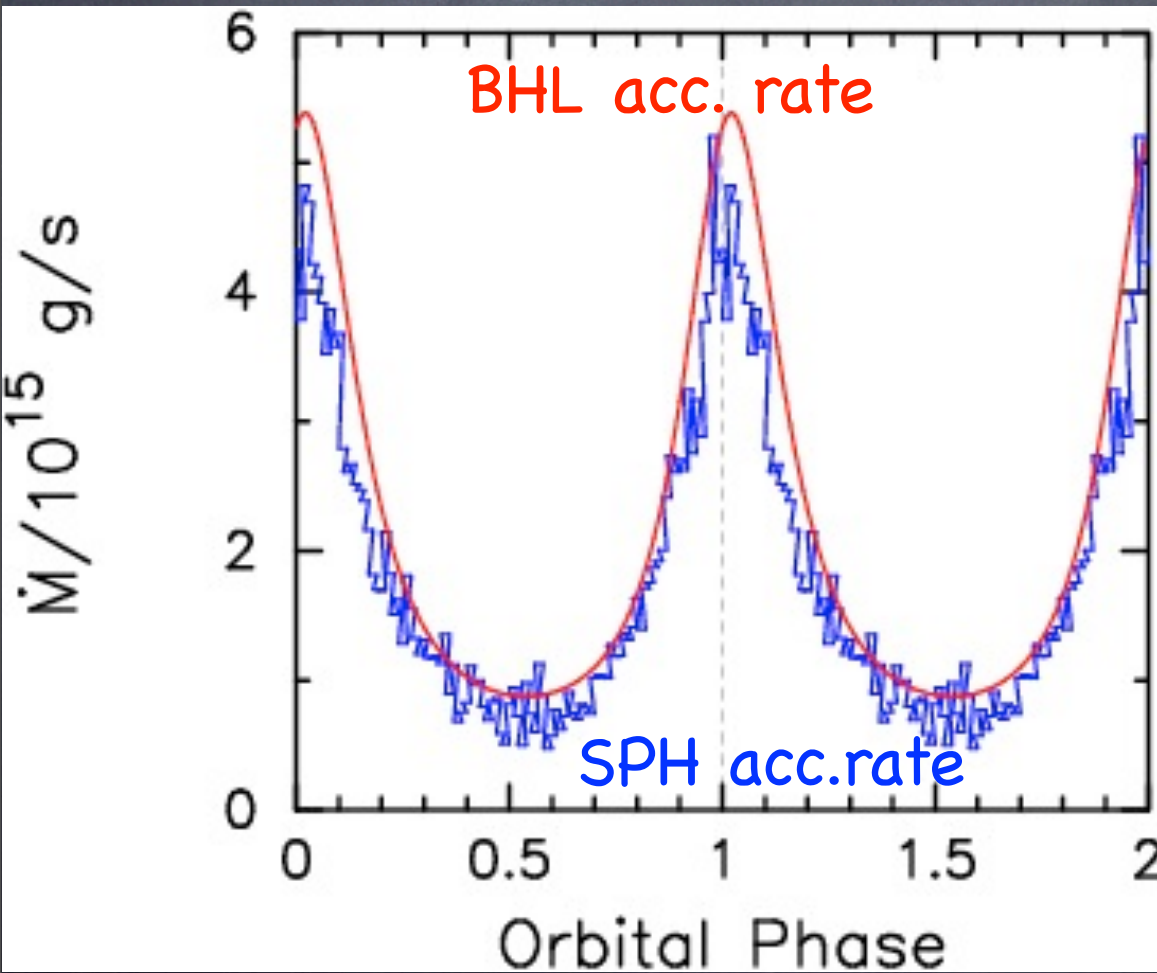
$$b = \frac{GM_{BH}}{V_{rel}^2 / 2}$$

BHL accretion rate

$$\dot{M}_{BHL} = \rho V_{rel} \pi b^2 = \frac{G^2 M_{BH}^2 \dot{M}_w}{V_{rel}^3 V_w d^2}$$



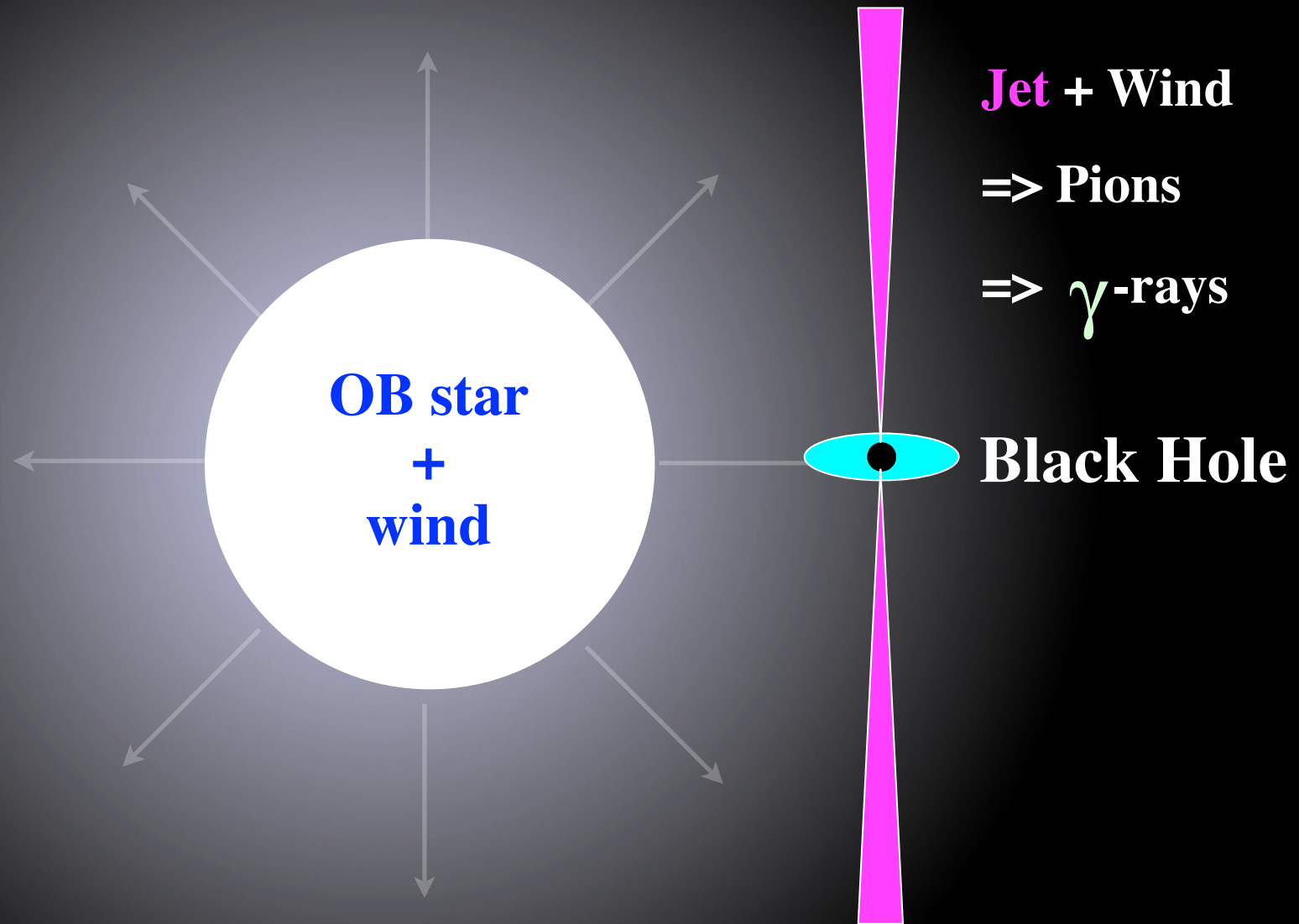
# Orbital variation of SPH accretion closely follows Bond-Hoyle-Lyttleton rate



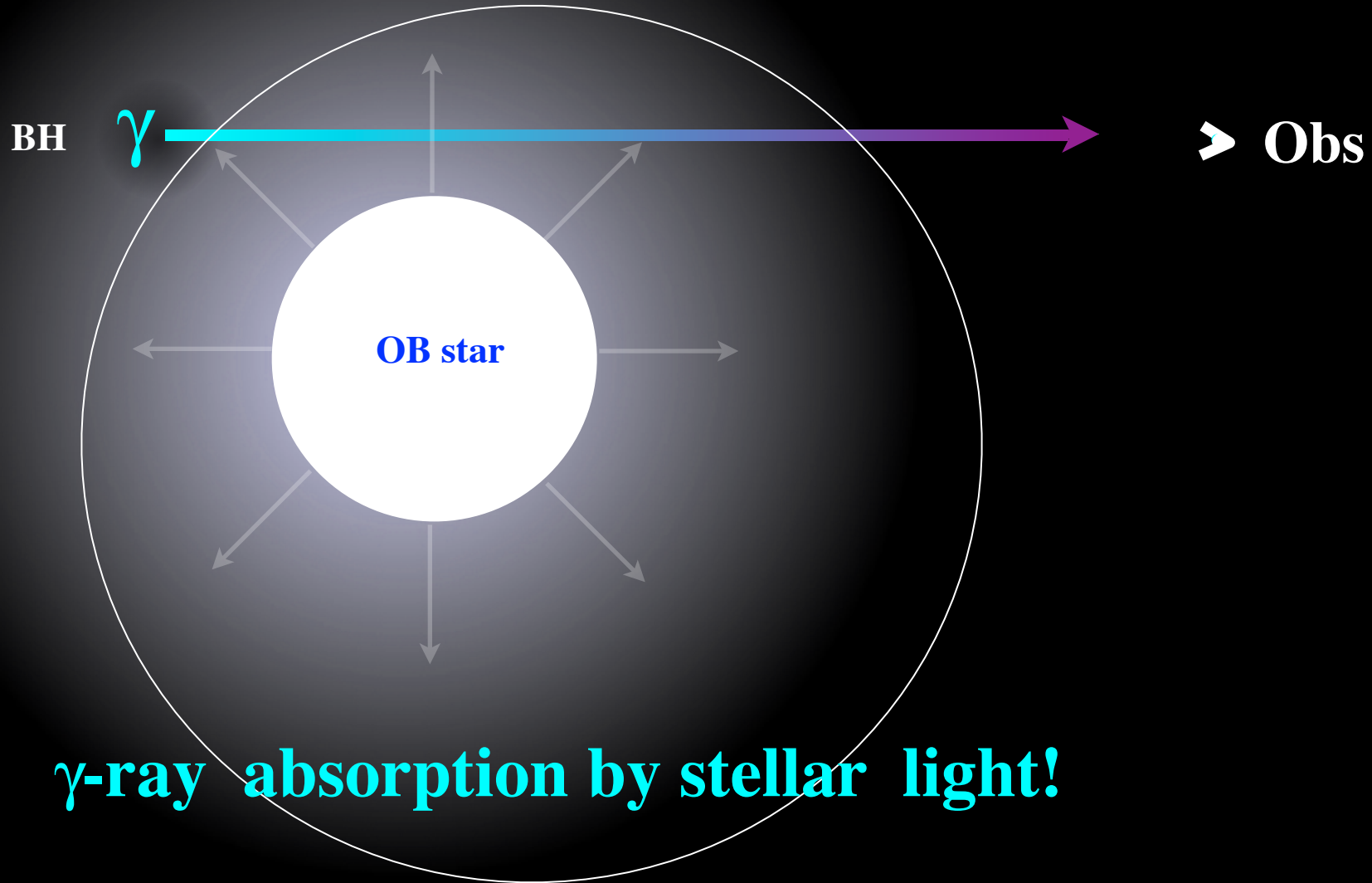
$$\dot{M}_{BHL} = \frac{G^2 M_X^2 \dot{M}_w}{V_{rel}^3 V_w d^2}$$

$$\beta = 1$$

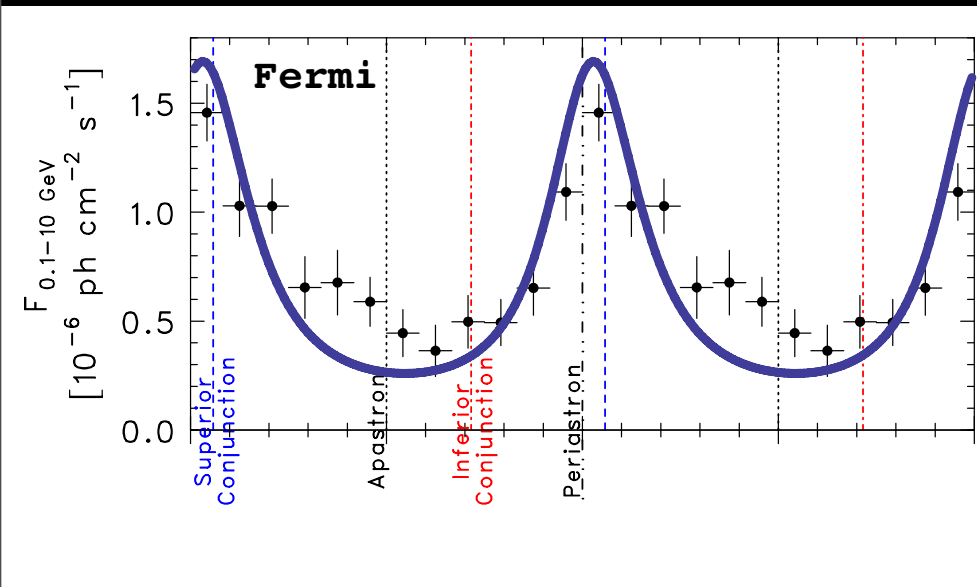
# MicroQuasar Model



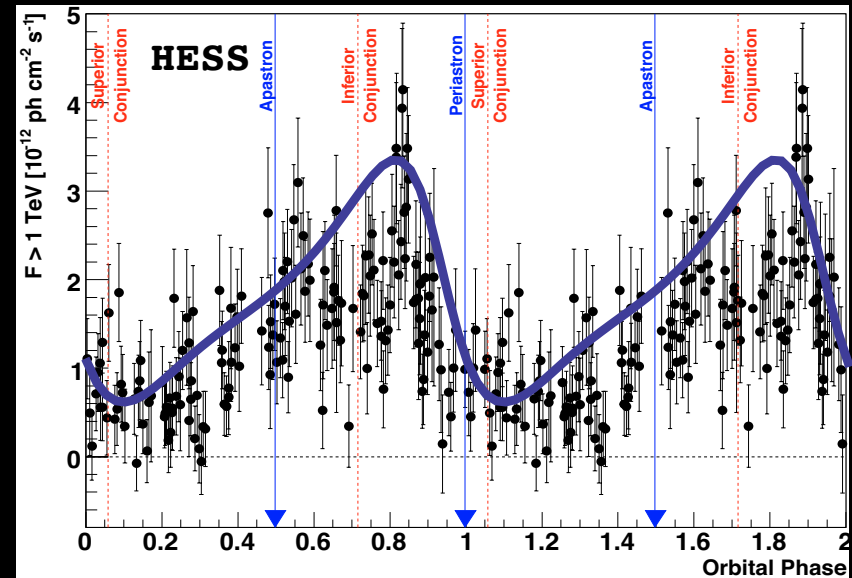
# $\gamma$ - $\gamma$ Absorption



# LS 5039 $\gamma$ -ray light curves

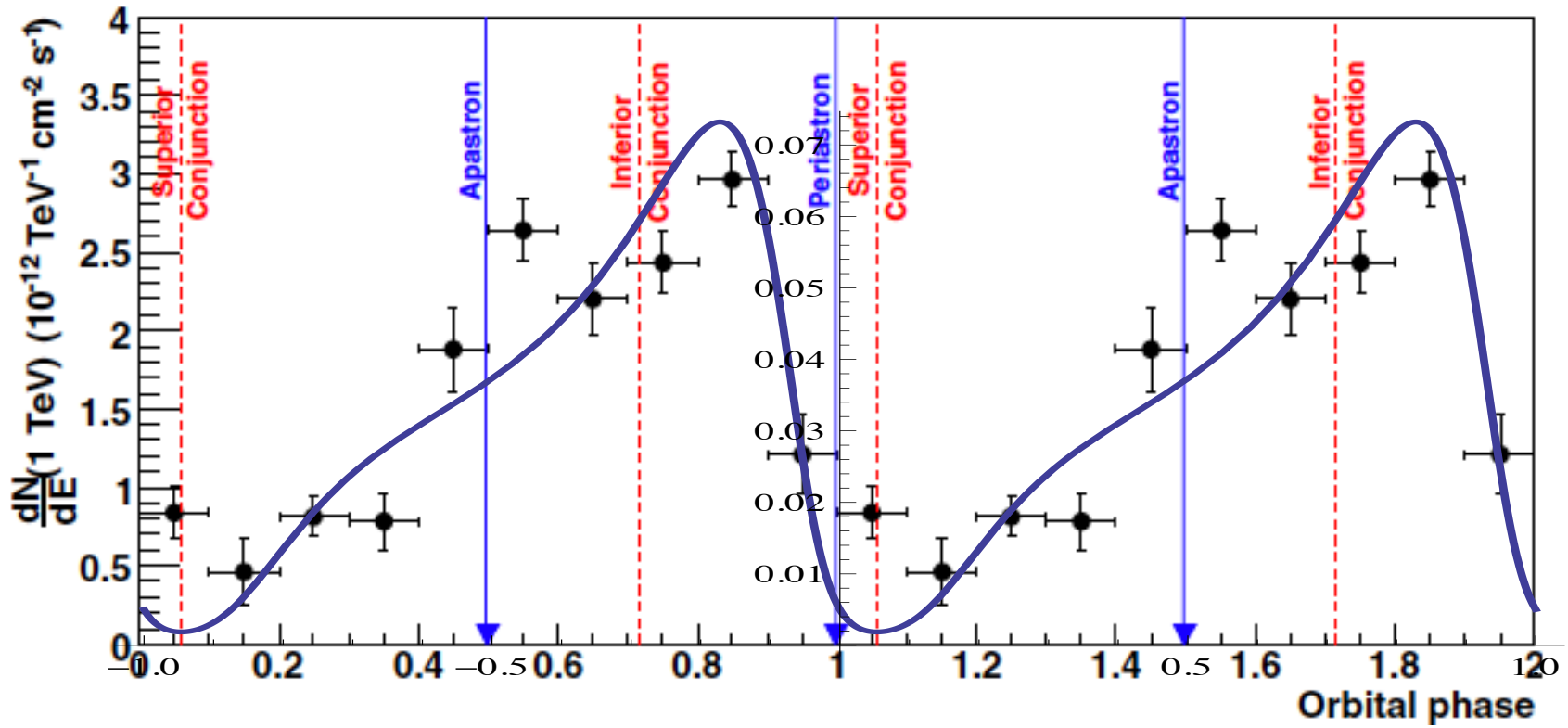


E=0.1-10 GeV:  
below threshold for  
 $\gamma$ - $\gamma$  with stellar UV  
**NO ABS**

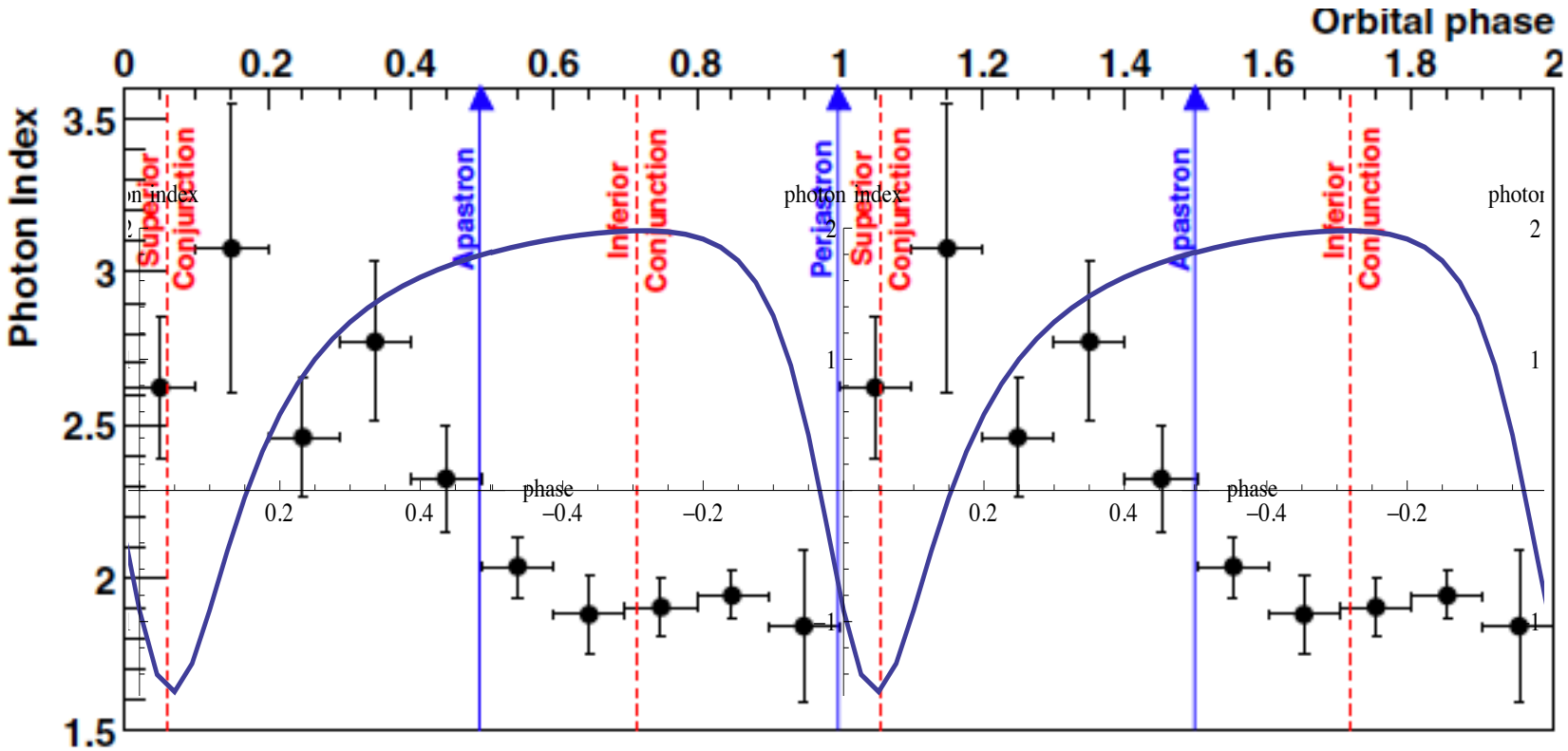


E > 1 TeV:  
above threshold for  
 $\gamma$ - $\gamma$  with stellar UV  
 **$\gamma$ - $\gamma$  ABS**

HESS  
normalization at E= 1 TeV.



# HESS photon index.



Need to include **Photon Cascade**

# Summary

- Winds unstable  $\Rightarrow$  clumps, soft X-rays
- Colliding Winds in Binaries  $\Rightarrow$  harder X-rays
- Be CSM = polar wind + eq. VDD + ablation
- Porosity length  $h = \text{size}/f_{\text{vol}}$  key clump parameter
- In HMXRB  $\Rightarrow$  fluctuations in X-ray light curve
- Microquasar jets  $\Rightarrow$   $\gamma$ -rays w/ fluct.  $\sim \text{Sqrt}[h/a]$
- MQ model fits Fermi & HESS l.c. for LS5039
- But fitting spectrum requires photon cascade

