

# Radiation Absorption and Reprocessing in $\gamma$ -Ray Binaries

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# Outline

- 1 Introduction
- 2 The gamma-ray binary scenario and secondary emission
- 3 Basic properties of the secondary emission
- 4 More detailed calculations
- 5 Summary

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# Gamma-ray binary systems: subclasses

- High-mass microquasars: massive star plus compact object that accretes forming jets.
- Pulsar high-mass binary: massive star plus pulsar with colliding winds.
- Massive star binary: two massive stars with colliding winds.

(e.g. Aharonian et al. 2005a; Albert et al. 2006; Acciari et al. 2008; Tavani et al. 2009a, 2009b; Abdo et al. 2009, 2010a)

- Low-mass WD binary (V407 Cygni)
- Low-mass (ms) pulsar binaries (J102347.6+003841)

(Abdo et al. 2010b; Tam et al. 2010)

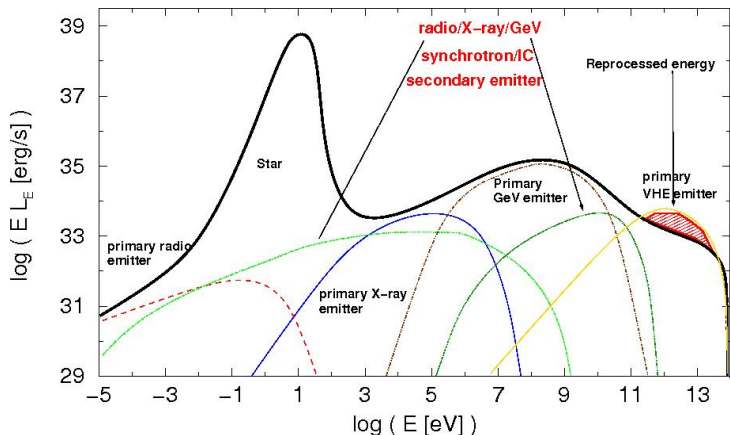
# Detected gamma-ray binaries with massive companions

- Three gamma-ray binaries have been detected in both the GeV and the TeV range: LS 5039, LS I +61 303 and PSR B1259–63.  
(Aharonian et al. 2005a, Albert et al. 2006, Acciari et al. 2008)
- Cygnus X-1 could be a TeV emitter as well...  
(Albert et al. 2007)
- Other candidates: HESS J0632+057, HD 215227 / AGL J2241+4454.  
(Hinton et al. 2009, Skilton et al. 2009, Falcone et al. 2010; Williams et al. 2010)
- Cygnus X-3,  $\eta$ Car and possibly Cygnus X-1 have been detected in GeV.  
(Tavani et al. 2009a,b; Abdo et al. 2009, 2010a; Sabatini et al. 2010)

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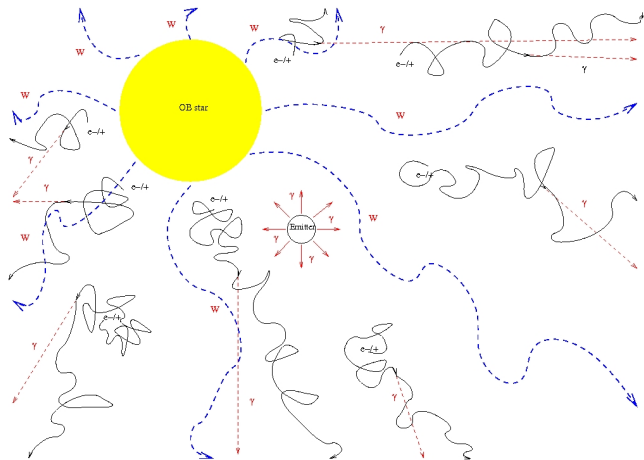
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# The non-thermal emission



Different components contribute to the non-thermal spectrum in gamma-ray binaries, of primary (accelerated particles) and secondary nature (gamma-ray absorption, hadronic process...).

# Sketching the physical system



Gamma-rays produce pairs in the system that get deflected and confined in the stellar wind, moving with it, radiating and cooling down.

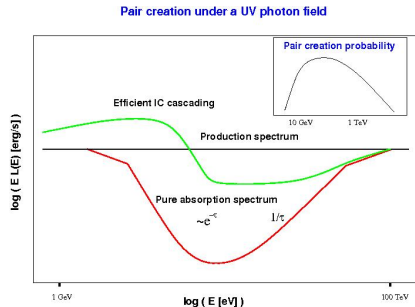


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# Pair creation

- Pair luminosity:  
 $\propto (1 - \exp^{-\tau}) \sim$   
 $(1 - \exp^{-0.5 L_{*38} d_{*12.5}^{-1}}) \sim 1$
- Energy threshold:  
 $E_{\text{th}} \sim 2/E_*(1 - \cos \theta) \sim 1/E_*$   
(pair distribution maximum)
- Interaction probability:  
 $\propto (1 - \cos \theta)$  (most of pairs  
toward the target source)



# Relevant timescales

- **Stellar IC scattering (Th.):**  $t_{\text{cool}} \geq 55 L_{*38}^{-1} E_{10 \text{ GeV}}^{-1} \text{ s}$
- **Synchrotron emission:**  $t_{\text{sync}} \sim 400 B_{10}^{-2} E_{10 \text{ GeV}}^{-1} \text{ s}$
- **Adiabatic cooling:**  $t_{\text{ad}} \sim 2 \times 10^4 d_{12.5*} \text{ s}$
- **Relativistic Bremsstrahlung:**  $t_{\text{br}} \sim 10^6 n_{w9}^{-1} \text{ s}$
- **Ionization cooling:**  $t_{\text{ion}} \sim 3 \times 10^7 E_{10 \text{ GeV}} n_{w9}^{-1} \text{ s}$
- **System light crossing time:**  $t_l = 100 d_{12.5} \text{ s}$
- **Diffusion time:**  $t_{\text{diff||}} \approx 10^{-3} E_{10 \text{ GeV}} \delta B_G^{-1} \text{ s}$

# Confinement/radiation regime

- $d_*/c < t_{\text{diff}\parallel} < t_{\text{rad}}$ : no isotropization, unefficient radiation
- $d_*/c < t_{\text{rad}} < t_{\text{diff}\parallel}$ : no isotropization, unefficient radiation
- $t_{\text{diff}\parallel} < d_*/c < t_{\text{rad}}$ : isotropization, adiabatic/radiation cooling
- $t_{\text{diff}\parallel} < t_{\text{rad}} < d_*/c$ : isotropization, efficient radiation
- $t_{\text{rad}} < d_*/c < t_{\text{diff}\parallel}$ : no isotropization, efficient radiation
- $t_{\text{rad}} < t_{\text{diff}\parallel} < d_*/c$ : no isotropization, efficient radiation

# Stellar wind magnetic field and particle confinement

- Beyond the Alfvén surface,  $\sim 1 - 2 R_*$ , the magnetic field becomes dynamically dominated by the wind material.
- At few  $R_*$ ,  $B_\phi \propto d_*^{-1}$  becomes dominant over  $B_r \propto d_*^{-2}$ .

(e.g. Usov & Melrose 1992)

- The condition of confinement,  $t_{\text{diff}\parallel} < d_*/c$  requires  $\delta B > 10^{-5} E_{10 \text{ G}} d_{*12.5} \text{ G}$ .
- Effective suppression of cascades requires  $B > 2 L_{*39}^{1/2} d_{*12.5}^{-1} E_{\text{TeV}}^{-0.85} \text{ G}$ .

# General properties

- Luminosities:

- $L_{e\pm, IC} \sim 5 \times 10^{-11} L_{\gamma 35} L_{*38} d_{*12.5}^{-1} \text{ erg cm}^{-2} \text{ s}^{-1}$
- $L_{\text{syncX}} \sim 10^{-11} L_{e\pm 35} B_{10G}^2 L_{*38}^{-1} \text{ erg cm}^{-2} \text{ s}^{-1}$
- $L_{5 \text{ GHz}} \sim 4 L_{e\pm 35} B_G \text{ mJy} \quad (2 \text{ kpc})$

- Spectrum (naively...):

- above  $E_{\text{th}}$ , flat (synch.) and steep (IC)
- below  $E_{\text{th}}$ ,  $\propto \nu^{-1.5}$  (synch./IC)

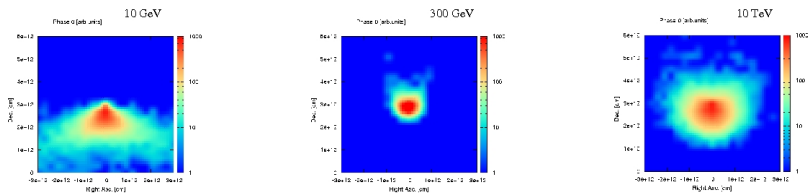
- Variability: geometry,  $L_{\gamma}(E_{\gamma})(\phi, t)$ ,  $B$ , wind inhomogeneities/anisotropies

- Morphology: spiral like/elongated at low energies, dependent on the eccentricity and the wind anisotropies/inhomogeneities, other outflows...

# Outline

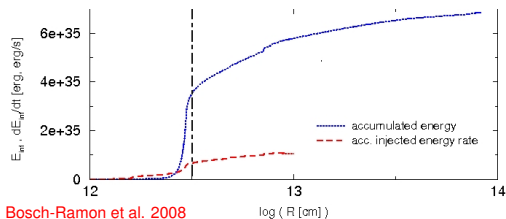
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# Spatial distribution of pairs



Bosch-Ramon & Khangulyan 2010, sub.

- Pairs distribute in the system depending on their energies.
- Significant amount of energy accumulate in the wind at scales of the binary system.

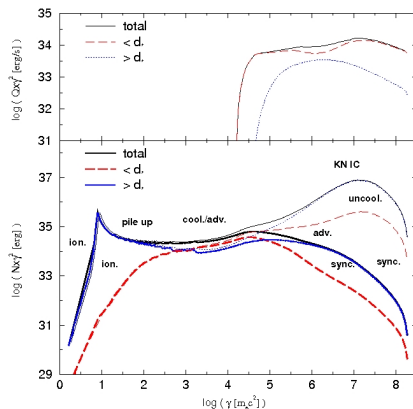


Bosch-Ramon et al. 2008



# Energy distribution of pairs

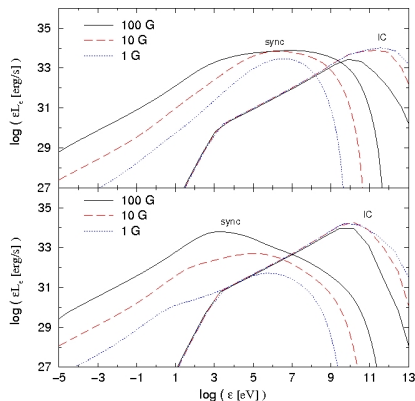
- $Q(E_{\text{th}}) \propto \sigma(E_{\gamma})_{\text{UV}} \rightarrow$ 
  - $N(E < E_{\text{th}}) \propto E^{-2, -1}$   
(Th. IC/synch. vs adiabatic)
  - $N(E > E_{\text{th}}) \propto Q(E) E^{+1, -1}$   
(KN IC vs synch.)
- Pairs from different regions have different spectra.



Bosch-Ramon et al. 2008

# The radiation spectral energy distribution

- For high  $B$  the synchrotron SED peaks at X-rays.
- The IC component peaks around 10 GeV, strongly softening at higher energies.
- For low  $B$ , KN IC cooling hardens all the spectra strongly  $\rightarrow$  IC cascade

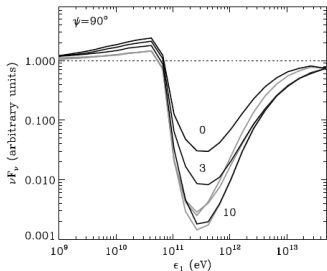


Secondaries can *violate* the synch.  $\nu_{\text{max}}$ -limit at  $\sim 100$  MeV.

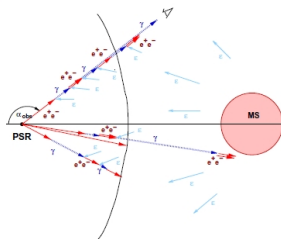
Bosch-Ramon et al. 2008

# Efficient cascading

- Very low  $B$  ( $B \ll 10^{-3} (E_{\text{TeV}}/d_{*12.5})$  G): 1 D
- Fairly low  $B$  ( $B \ll 1 (L_{*39}/d_{*12.5})$  G): pure 3 D
- Intermediate cases:  $L_{\text{TeV}} \propto B^{-2}$



(Cerutti et al. 2010)

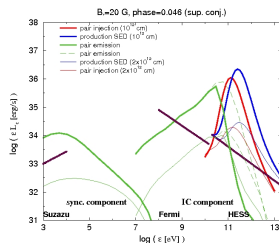


(Sierpowska-Bartosik & Torres 2008)

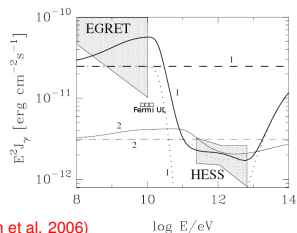
(see also Aharonian et al. 2006; Orellana et al. 2007; Khangulyan et al. 2008; Bosch-Ramon et al. 2008; Bednarek 2010)

# The case of LS 5039 (I)

- Efficient gamma-ray production inside the binary system can easily violate the X-ray and GeV fluxes.
- A pure 3D IC cascade model cannot explain the data above GeV in LS 5039.
- Photon-photon absorption cannot play a dominant role → system periphery emitter



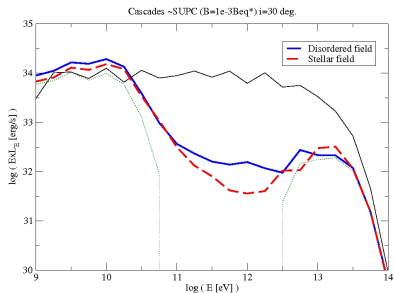
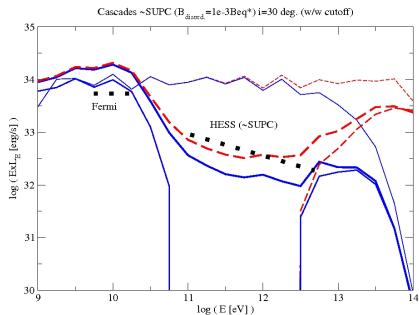
preliminary



(Aharonian et al. 2006)

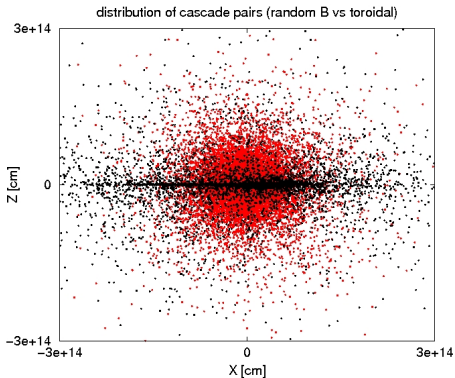
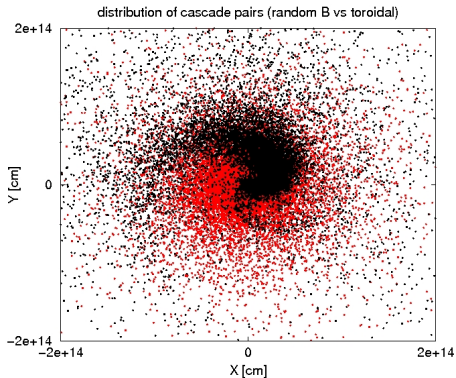
# The case of LS 5039 (II)

Another exemple of a 3D cascade (low  $B$ ) at SUPC in LS 5039



preliminary

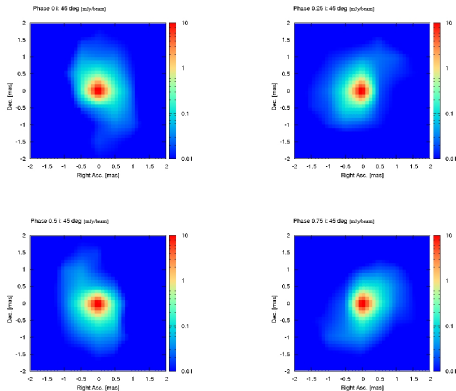
# Pair distribution: stellar vs disordered $B$



Diffusion spreads up/downwards, but confines sideways.

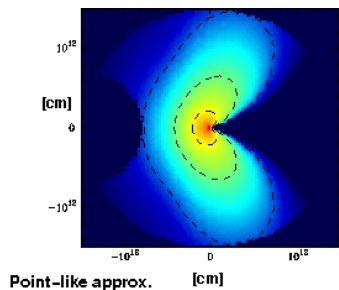
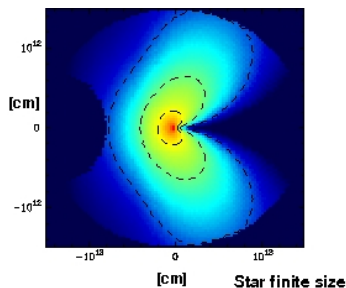
# Radio emission from secondary pairs in gamma-ray binaries

5 GHz emission in the observer plane produced by secondary pairs created in a gamma-ray binary.



## Pair injection: point-like versus finite size

For secondary radio emission the star point-like approximation is fairly good.



LS 5039 : ~ periastron

Bosch-Ramon & Khangulyan 2010, sub.



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# Summary

- Gamma-rays with  $E_\gamma > 10$  GeV fill the binary system.
- Pairs are created under the presence of a dense (stellar) photon field.
- Confinement of particles is likely efficient (isotropic pair distribution, 3D cascades).
- For moderate magnetic fields:
  - Cascades get suppressed
  - X-ray emission can overcome the observed fluxes  $\rightarrow B_x \lesssim 10$  G.
- In LS 5039, absorption nor cascades should be dominant.

Open questions: localization/extension of the emitter,  $B$  geometry,  $B$  in radio regions, impact on the wind...